

# **Chapter II**

**Synthesis, characterization and growth mechanism  
of Gd<sub>2</sub>O<sub>3</sub> nanosystems**

## 2.1 Principle and characterization techniques employed

Synthesis of nanomaterial is a topic of intensive investigation in both basic and applied research. Large scale production, cost effectiveness and stability of nanoscale systems are among key issues for their deployment in specific applications. In this regard, size and shape controlled nanosystems can be acquired through suitable fabrication techniques. Essentially, an appropriate growth mechanism is the key to yield quality nanomaterial products while introducing the desired phase, shape and also narrow size distribution. The present chapter focuses on different approaches adopted for synthesizing undoped and doped Gd<sub>2</sub>O<sub>3</sub> nanoscale structures. The structural, optical and morphological evolutions are also highlighted.

Pure and doped Gd<sub>2</sub>O<sub>3</sub> nanosystems have been fabricated via physicochemical and hydrothermal routes. The crystallographic information was revealed through a Rigaku miniFlex X-ray diffractometer (XRD) that employs a  $CuK_{\alpha}$  source ( $\lambda=1.543 \text{ \AA}$ ). The diffraction angle ( $2\theta$ ) was varied in the range of  $20^{\circ}$ - $60^{\circ}$ , and in steps of  $0.05^{\circ}$ . Transmission electron microscopy (TEM) images were obtained either from a JEOL JEM 2100 or a TECNAI G2 20 S-TWIN machine, operating at an accelerating voltage of 200 kV. The selected area electron diffraction (SAED) pattern also performed with the help of TEM. The room temperature electron paramagnetic resonance (EPR) responses have been acquired on a JEOL: JESFA200 EPR spectrometer and considering a frequency 9.37 GHz and field modulation 100 kHz. The optical absorption spectra have been recorded on a UV-Visible spectrophotometer (Model UV2450, Shimadzu Corporation, Japan) working in the range of wavelength 200-700 nm. The room temperature PL spectra of the samples were revealed by using a LS 55 PerkinElmer spectrophotometer. The room temperature vibrational features were explored by a micro-Raman spectrometer (Renishaw, UK) and considering an excitation wavelength 514.5 nm of Ar<sup>+</sup> line. Moreover, for magnetic measurements were carried out on an SQUID (Evercool SQUID VSM DC magnetometer, Quantum design, USA).

## 2.2 Gd<sub>2</sub>O<sub>3</sub> nanoparticles

### 2.2.1 Synthesis and characterization of undoped and Tb<sup>3+</sup> doped Gd<sub>2</sub>O<sub>3</sub> nanoparticles

First, gadolinium acetate hydrate [Gd(CH<sub>3</sub>CO<sub>2</sub>)<sub>3</sub>·H<sub>2</sub>O] (GdAc, CDH, 99.9% pure) is dispersed in deionized water and then a 1:1 molar sol has been prepared with citric acid [99% pure, CDH] under vigorous stirring for 30 min. The mixture is placed in an open-air oven at a temperature of 70°C in order to facilitate chemical reduction of GdAc to Gd(OH)<sub>3</sub> and to remove away the undesired water content. The product was further sintered in a tubular furnace and at a temperature of 400°C for 4 h. Later, the product is calcined at a temperature of 800°C for 2 h so as to discard byproducts and carbon-rich compounds. For synthesizing Tb<sup>3+</sup> doped Gd<sub>2</sub>O<sub>3</sub> product, terbium acetate Tb(CH<sub>3</sub>CO<sub>2</sub>)<sub>3</sub>·H<sub>2</sub>O] (TbAc, CDH, 99.9% pure) is added to the reaction mixture in the initial stage of the reaction, whereas rest of the reaction steps were kept similar as mentioned above. Quantitatively, in order to make 0.5, 1, 3 and 5% Tb<sup>3+</sup> doped Gd<sub>2</sub>O<sub>3</sub> powders, the Gd to Tb molar ratio was varied as 0.005, 0.01, 0.03 and 0.05; respectively. The flow chart of the overall procedure is illustrated in Figure 2.1.

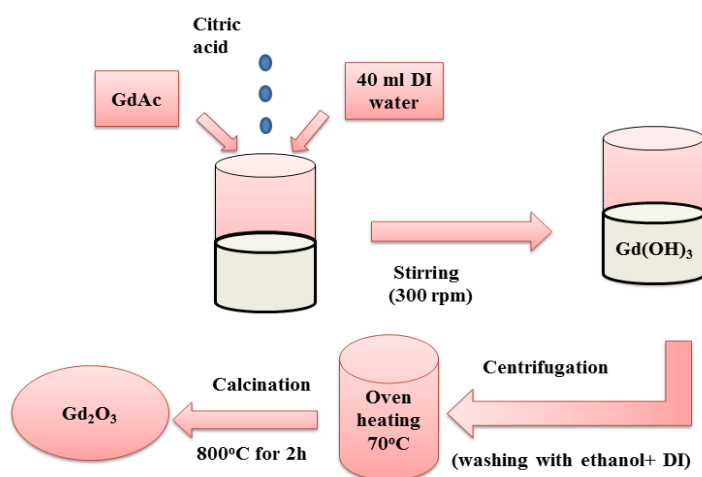


Figure 2.1: Synthesis protocol for preparing Tb<sup>3+</sup> doped Gd<sub>2</sub>O<sub>3</sub> nanosystems

Figure 2.2 (a) shows a series of the XRD patterns, of undoped and Tb<sup>3+</sup> doped Gd<sub>2</sub>O<sub>3</sub> nanoparticles. The four prominent peaks, located at 28.50, 33.10, 47.60, and 56.55° corresponded to (222), (400), (440) and (622) crystallographic planes. With a preferred orientation along (222) plane, the patterns characterize cubic crystalline phase of Gd<sub>2</sub>O<sub>3</sub> (JCPDS No.11-0604) [1]. As no appreciable peak shifting was noticeable in case of Tb-doped Gd<sub>2</sub>O<sub>3</sub> system, the Tb<sup>3+</sup> ions are expected to have undergone substitutional doping at Gd<sup>3+</sup> sites [1].

The average crystallite size (*d*) and microstrain (*η*) can be calculated from the Williamson-Hall (W-H) expression (Appendix III) [2]:

$$\beta \cos \theta = 0.9 \lambda / d + 4 \eta \sin \theta, \quad (2.1)$$

where  $\beta$  is the full width at half maxima,  $2\theta$  is the Bragg diffraction angle in degrees,  $\lambda$  is the wavelength of X-ray (1.543 Å). From W-H plots, the average crystallite size was found to be in the range of 11 - 24 nm and with the average size nearly doubled for the doped system as compared to its undoped counterpart (Table 2.1). The microstrains are of the order of 10<sup>-3</sup> and 10<sup>-4</sup>. While exhibiting a negative microstrain (10<sup>-3</sup>), the undoped Gd<sub>2</sub>O<sub>3</sub> nanopowder was believed to be comprised of much relaxed crystallites. With Tb doping, however, microstrain becomes positive (10<sup>-4</sup>). Though the magnitude is ~10 times smaller than the undoped case, a reasonable lattice mismatch and consequently, internal stress might have occurred due to inclusion of Tb<sup>3+</sup>. The magnitude of microstrain is doubled when the doping level reached ~5%.

The morphological feature of the nanoparticles is revealed through HRTEM imaging, depicted in Figure 2.2 (b). The size distribution of the nanoparticles systems is shown as figure-inset, in the form of a histogram plots. The average diameter of the nearly spherical nanoparticles is found to be approximately 30 nm.

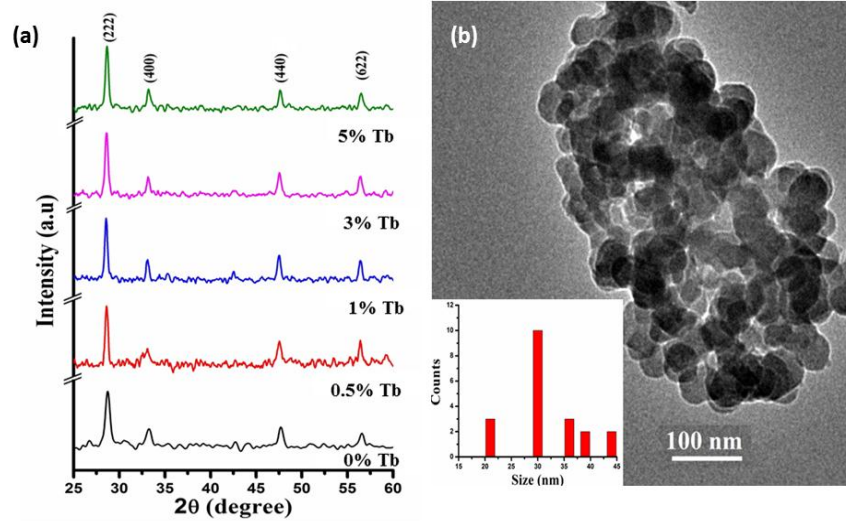


Figure 2.2: (a) A series of X-ray diffractograms of undoped and Tb<sup>3+</sup> doped Gd<sub>2</sub>O<sub>3</sub> nanophosphors, (b) HRTEM micrograph of the nanoscale Gd<sub>2</sub>O<sub>3</sub> system.

Table 2. 1: Properties of Tb<sup>3+</sup> doped Gd<sub>2</sub>O<sub>3</sub> nanocrystalline phosphors

Tb concentration (%)	Avg. crystallite size ( <i>d</i> ) (nm)	Microstrain ( $\epsilon$ )
0	11.8±2.1	-2.15×10 <sup>-3</sup>
0.5	21.6±2.3	2.35×10 <sup>-4</sup>
1.0	24.7±1.2	2.18×10 <sup>-4</sup>
3.0	24.7±1.5	1.98×10 <sup>-4</sup>
5.0	24.4±1.2	4.32×10 <sup>-4</sup>

## 2.2.2 Optical characteristics of undoped and Tb<sup>3+</sup> doped Gd<sub>2</sub>O<sub>3</sub> nanoscale systems

Optical absorption spectra of the un-doped and Tb<sup>3+</sup> -doped Gd<sub>2</sub>O<sub>3</sub> nanoparticles are shown in Figure 2.3. The samples show three characteristic peak maxima located at ~ 254 nm, 274 nm and 350 nm. Know that, the energy absorption in Gd is mediated via transition of 4*f* electrons to 5*d* level and then reorganization of the 4*f* electrons into various multiplets [3]. The seven electrons in the 4*f* orbitals of Gd have as many as 3432 multiplets and the ground state is <sup>8</sup>S<sub>7/2</sub> [3]. Among these multiplets, the low lying multiplet above the ground state is <sup>6</sup>P<sub>7/2</sub>. The absorption peak at ~274 nm is typical for Gd<sub>2</sub>O<sub>3</sub>

particles and is attributed to  $^8S_{7/2} \rightarrow ^6I_{7/2}$  transition. Possibly, the weak peak at  $\sim 254$  nm is due to one of the transition corresponding to  $^8S_{7/2} \rightarrow ^6I_{j/2}$  multiplets where  $j=9, 12, 13, 15$  and  $17$  [4]. The peak at  $\sim 350$  nm is ascribed to the  $^6P_{7/2}-^8S_{7/2}$  of Gd(III) [5].

With the symbols having their usual meaning, and using generalized Tauc relation [6]:  $(\alpha h\nu)^{1/n} = A(h\nu - E_g)$ , one can estimate optical band gap of the systems under investigation. Here  $h\nu$  is the incident photon energy,  $\alpha$  is the absorption coefficient,  $E_g$  is the optical band gap, and  $A$  is the system dependent parameter. Assuming all transitions are direct, and allowed type ( $n=1/2$ ), the estimated  $E_g$  values was found to decline with increase in Tb<sup>3+</sup> doping. As for the undoped Gd<sub>2</sub>O<sub>3</sub> phosphor,  $E_g = 5.22$  eV and in consistency with an earlier report [7]. For 0.5, 1, 3 and 5% Tb doped systems, the gap was estimated as  $\sim 5.18, 5.17, 5.16$  and  $5.14$ ; respectively. The reduction of the fundamental band gap with increasing doping concentration suggests creation of intermediate states below the bottom of the conduction band. The localized states could either be due to the presence of adequate oxygen vacancies or/and impurity states due to Tb-incorporation. Also, band gap follows a decreasing trend with the increase in size of the crystallites [8].

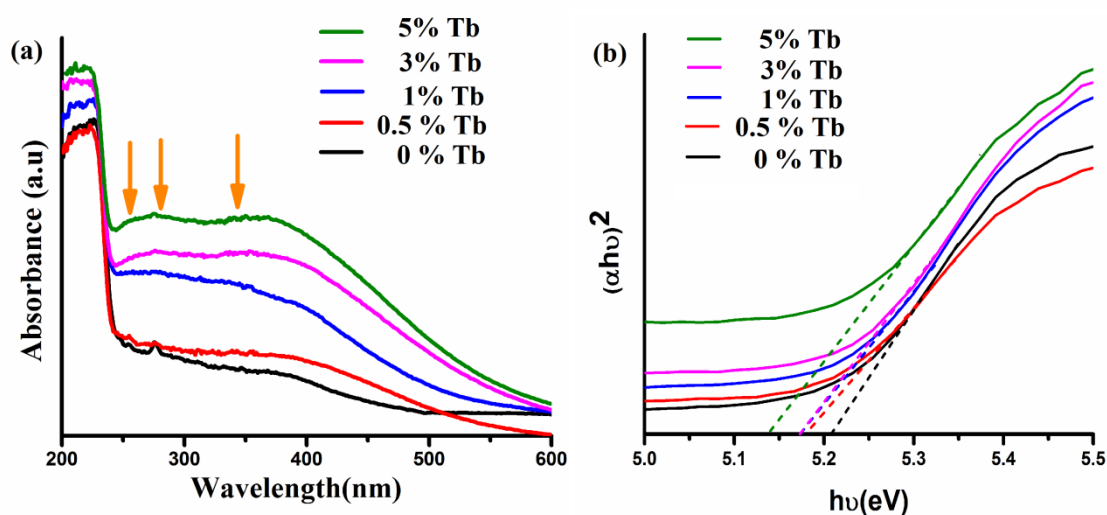


Figure 2.3: (a) Optical absorption spectra of undoped and Tb<sup>3+</sup> doped Gd<sub>2</sub>O<sub>3</sub> nanophosphors. Tauc plots of respective samples are shown in (b).

Furthermore, we have calculated the refractive indices ( $\mu$ ) of the nanophosphors using popular Dimitrov formula [9]:

$$\mu^2 - 1/\mu^2 + 2 = 1 - \sqrt{(E_g/20)} \quad (2.2)$$

For the undoped Gd<sub>2</sub>O<sub>3</sub> phosphor,  $\mu = 1.97$  [10]. It has increased to 1.976, 1.977, 1.979, and 1.982 for 0.5, 1, 3 and 5% doped Gd<sub>2</sub>O<sub>3</sub> systems. In other words, refractive index is improved by  $\sim 0.6\%$  for a maximum Tb<sup>3+</sup> doping (of 5%). Accordingly, the doped systems are optically more dense as compared to the undoped Gd<sub>2</sub>O<sub>3</sub> counterpart. Thus the material becomes optically more active in view of inclusion of additional localized states into the system.

The PL emission spectra of the undoped and Tb<sup>3+</sup> doped Gd<sub>2</sub>O<sub>3</sub> nanophosphors (under 300 nm excitation) are depicted in Figure 2.4. In a typical spectrum, three main emission peaks were found to be located at  $\sim 432$  nm, 450 nm and 492 nm but with variable emission intensities. The peak positions have been extracted through normalized multipeak Gaussian fitting such that the area under the experimental curve and the empirical curve remains same. The peak maxima at  $\sim 432$  and 450 nm can be attributed to the surface defects of Schottky and Frankel kinds, whereas the peak maxima at  $\sim 492$  nm are attributed to the self-trapped exciton luminescence [11]. Note that, in case of Tb<sup>3+</sup> doped Gd<sub>2</sub>O<sub>3</sub> system, the emission peak at  $\sim 492$  nm is adequately improved owing to superimposition of  $^5D_4 \rightarrow ^7F_6$  related transitions of Tb<sup>3+</sup>, which occur almost in the same wavelength regime. However, the transitional events are rather low for a higher level of Tb<sup>3+</sup> doping. Interestingly, the broad emission peak at  $\sim 548$  nm is most intense for the 0.5% Tb<sup>3+</sup> doped Gd<sub>2</sub>O<sub>3</sub> system as compared to other samples. The green emission due to the Tb<sup>3+</sup> mediated  $^5D_4 \rightarrow ^7F_5$  transition, which is absent in case of the undoped system, was found to get quenched for a higher amount Tb<sup>3+</sup> inclusion. There is a possibility that some interstitial sites can be occupied at higher doping levels of Tb<sup>3+</sup>. We speculated that this transition is mediated by Tb<sup>3+</sup> ions present in the interstitial sites. Since diffraction technique differs from spectroscopy technique this could not be ascertained.

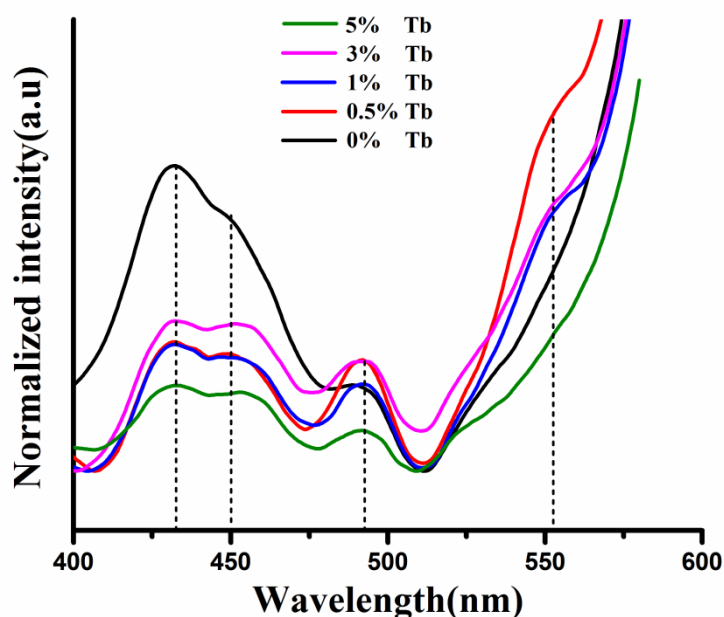


Figure 2.4: Room temperature photoluminescence spectra of the undoped and Tb<sup>3+</sup> doped Gd<sub>2</sub>O<sub>3</sub> nanophosphors ( $\lambda_{ex}$  = 300 nm).

The PL intensity ratio between  ${}^5D_4 \rightarrow {}^7F_5$  and  ${}^5D_4 \rightarrow {}^7F_6$  transitions was estimated to be 3.8 for 0.5% Tb doping case, whereas it is approximately 2.4 for 1% and 3% Tb-doping cases. Thus it is clear that, only an increment of Tb<sup>3+</sup>-doping level cannot induce intense green emission responses. Although in a solvent media, Eu<sup>3+</sup> or Tb<sup>3+</sup> doped nanometric Gd<sub>2</sub>O<sub>3</sub> gives a maximum emission signature corresponding to a dopant concentration of 5–6% [12], we predict that solid Gd<sub>2</sub>O<sub>3</sub> powder could act as an efficient green phosphorescent candidate when the doping level is ~0.5%. This is because, the presence of hydroxyl groups, present on the nanocrystallites, can quench the emission response significantly. It is also worth mentioning here that, for a high doping level one may have to sacrifice some amount of crystallinity in the desired system.

The room temperature time resolved photoluminescence (TR-PL) data of 1% and 5% Tb<sup>3+</sup> doped Gd<sub>2</sub>O<sub>3</sub> nanosystem measured for emission wavelengths value  $\lambda_{em}$  = 490, and  $\lambda_{em}$  = 545 nm are shown in Figure 2.5. The decay curve followed a bi-exponential function, given by equation [12]



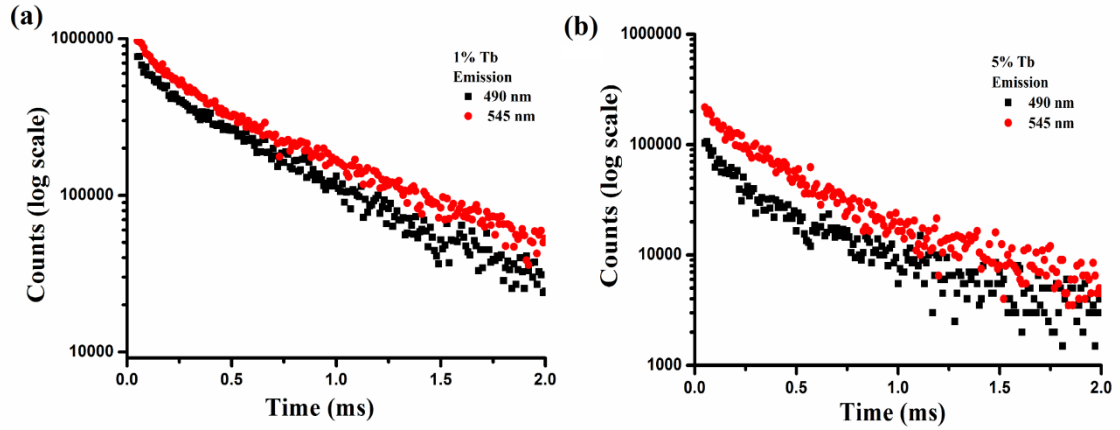


Figure 2.5: TR-PL spectra of (a) 1% and (b) 5% Tb doped Gd<sub>2</sub>O<sub>3</sub> nanoparticles. Black and red experimental curves represent the nature of PL decays corresponding to  $\lambda_{em} = 490$  nm and 545 nm, respectively.

$$A = A_1 e^{-t/\tau_1} + A_2 e^{-t/\tau_2} \quad (2.3)$$

where  $\tau_1$  and  $\tau_2$  are the decay time constants associated with the fast and slow processes and  $A_1$  and  $A_2$  as corresponding emission intensities.

The time constants associated with the 490 nm and 545 nm emissions are shown in Table 2.2. The decay life time of our samples is much lower than the predicted ones in the earlier report [12]. This effect is may be due to the fast energy dissipation by the existence of numerous non-radiative recombination channels. Moreover, the presence of structural defects and hydroxyl groups could also lead to quenching in the emission profile thereby lowering the decay time [13].

Table 2.2 : Decay time constants calculated from TR-PL spectra

Sample	Emission wavelength (nm)	$\tau_1$ (ms)	$\tau_2$ (ms)
1% Tb	490	0.58	0.07
	545	0.11	0.64
5% Tb	490	0.49	0.93
	545	0.04	0.36

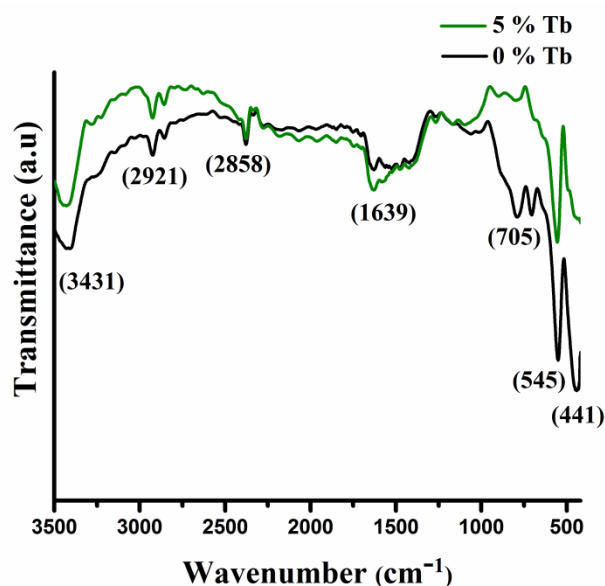


Figure 2.6: FTIR spectra of the undoped and Tb<sup>3+</sup> doped Gd<sub>2</sub>O<sub>3</sub> nanosystems.

Fourier transform infrared (FTIR) spectra of the undoped and 5% Tb<sup>3+</sup> doped Gd<sub>2</sub>O<sub>3</sub> phosphors are presented in Figure 2.6. The main bands located at ~ 2858 cm<sup>-1</sup> and 2921 cm<sup>-1</sup> can be assigned to the symmetric, asymmetric stretching vibrations of -CH<sub>2</sub> [1]. The band at ~705 cm<sup>-1</sup> is the C-O vibration of the synthesized sample [14]. The bands at ~3431 cm<sup>-1</sup> and 1639 cm<sup>-1</sup> are attributed to the O-H stretching vibration and bending of water molecules. Whereas, the bands at located ~545 cm<sup>-1</sup> and ~441 cm<sup>-1</sup> are caused by the stretching vibration of Gd-O of cubic Gd<sub>2</sub>O<sub>3</sub> [15]. It may be noted that, while other vibrational modes are not substantially modified, the strength of Gd-O vibrational stretching is adequately improved for the Tb<sup>3+</sup> doped Gd<sub>2</sub>O<sub>3</sub> nanosystem.

## 2.3 Gd<sub>2</sub>O<sub>3</sub> nanorods

### 2.3.1 Synthesis and characterization of Gd<sub>2</sub>O<sub>3</sub> nanorod powders

A simple hydrothermal procedure is carried out to grow quality Gd<sub>2</sub>O<sub>3</sub> nanorods. In a typical procedure, 0.28 g of bulk Gd<sub>2</sub>O<sub>3</sub> (Otto, 99.9% purity) is first dispersed in 40 ml distilled water followed by dropwise addition of dil. HNO<sub>3</sub> until a clear sol is obtained. The pH of the mixture is then adjusted to 9.2, 10.8, and 13.3 by adding several drops of 5M NaOH solution under vigorous stirring (~300 rpm). The mixture is then transferred to a teflon-lined stainless

steel autoclave and subjected to oven heating at a temperature of  $\sim 140^\circ\text{C}$ , for 24 h. A white solid product of Gd(OH)<sub>3</sub> is now collected by filtering the precursor using a whatman filter®. In the next step, the product was washed several times, first with distilled water and then with ethanol, followed by oven-drying at a temperature of  $80^\circ\text{C}$ . Finally, the hydroxide powder is annealed at  $600^\circ\text{C}$ , for 3h. The spontaneous decomposition and subsequent dehydration of Gd(OH)<sub>3</sub> is believed to be facilitated, yielding Gd<sub>2</sub>O<sub>3</sub> nanorod product. The chemical reaction steps and flow chart diagram of hydrothermally synthesized Gd<sub>2</sub>O<sub>3</sub> nanorods are as described below .

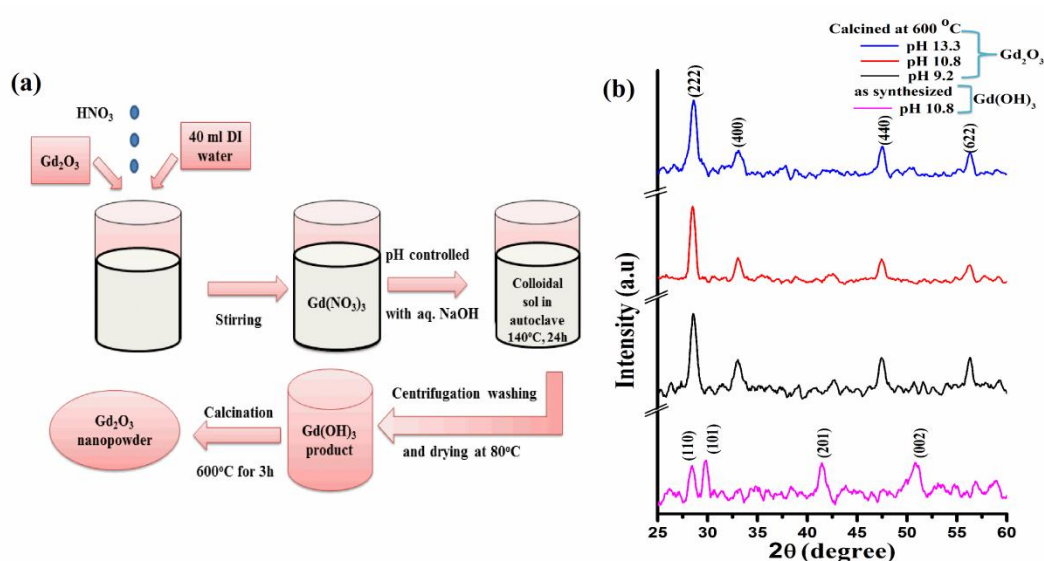
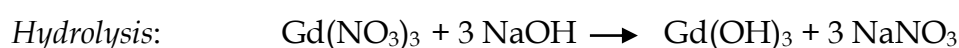
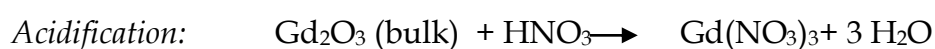


Figure 2.7: (a) Block diagram of hydrothermally synthesized Gd<sub>2</sub>O<sub>3</sub> nanorods, (b) XRD patterns of as-synthesized Gd(OH)<sub>3</sub> product (pH 10.8) and Gd<sub>2</sub>O<sub>3</sub> nanorods derived from precursors of different pH control.

The XRD patterns of as-synthesized Gd(OH)<sub>3</sub> product prepared at pH = 10.8 and Gd<sub>2</sub>O<sub>3</sub> nanorod-powders prepared at pH = 9.2, 10.8 and 13.3 are depicted in Figure 2.7(b). The hydroxide product exhibited a hexagonal crystal structure with observable diffraction peaks indexed as (1 1 1), (1 0 1), (2 0 1) and (0 0 2) crystallographics planes (JCPDS no. 83-2037) [16]. After sintering at a

temperature 600°C, prominent, yet apparently broadened diffraction peaks corresponding to (2 2 2), (4 0 0), (4 4 0) and (6 2 2) crystallographic planes are obtained ensuring the existence of a cubic crystalline phase (JCPDS no. 11-0604) [1]. No peak due either to hydroxides, impurities or byproducts has been detected.

With symbols having their usual meanings and using diffraction formula relevant to cubic phase:

$$d_{hkl} = (h^2+k^2+l^2/a^2)^{1/2} = \lambda/2 \sin \theta \quad (2.4)$$

The lattice parameter ( $a=b=c$ ) of the samples synthesized at pH of 9.2, 10.8 and 13.3, are estimated to be 10.80, 10.83 and 10.79 Å, respectively. Earlier, a cubic-phase cell parameter of magnitude 10.81 Å has been predicted for the Gd<sub>2</sub>O<sub>3</sub> system [17]. The average crystallite size ( $d$ ) and microstrain ( $\eta$ ) can be calculated by using popular W–H expression as mentioned in equation (2.1). The average crystallite sizes are found to be ~ 5.19, 6.46 and 6.09 nm as for pH 9.2, 10.8 and 13.3 cases, respectively (Table 2.3).

Table 2. 3. Average crystallite size (nm), lattice parameter (nm), microstrain and dimension of nanorods (nm) of as-synthesized Gd<sub>2</sub>O<sub>3</sub> nanorod systems.

Nanorods synthesized from precursors	Avg. crystallite size ( $d$ ) (nm)	Lattice parameter ( $a=b=c$ ) (nm)	Microstrain ( $\times 10^{-3}$ )	Dimension of nanorods		
				Length (nm)	Diameter (nm)	Aspect ratio
*pH 9.2	5.19	1.080	-12.7	---	---	---
pH 10.8	6.46	1.083	-8.9	75-115	19	~4
pH 13.3	6.09	1.079	-8.5	60-128	7	~8.5

\*TEM analysis not performed on this sample.

The HRTEM images of the nanorod samples, prepared at pH 10.8 and pH 13.3, are shown in Figure 2.8 (a) and (b). As for pH 10.8 specimen, a number of nanorods with an average diameter ~19 nm and length ranging from 76 to 115 nm can be evident in the micrograph (Figure 2.8(a)). With a smoothed surface, the rods resemble the structure of ‘unfold fingers’. Note that, for any nanorod, the diameter at the middle of the rod is much larger than at the edge. This is ascribed to the spontaneous decomposition of gadolinium hydroxide

into its oxide product followed by subsequent assimilation of crystallites along (222) plane and instant termination at the end. Also, all the Gd<sub>2</sub>O<sub>3</sub> nanorods were not found in the same focal plane and that is why we observed a reasonable difference in the image contrast of many off-aligned rods. As for sample synthesized from a precursor of pH 13.3, the range of average nanorod length was found to be wider i.e., 60-128 nm but having a lower average diameter  $\sim 7.1$  nm (Figure 2.8 (b)). Thus, both shorter and longer rods could be achievable at a higher pH. We believe that, unrestricted preferential growth of nanocrystallites might have led to longer sized rods. But shorter nanorods could have been developed due to obstruction of one nanorod growing in the vicinity of the other or, thermodynamic instability that would keep surface energy of the overall system minimum. Consequently, the diameter of the rods can also be reduced in this case. As diameter is close to the average crystallite dimension (see, XRD analysis) we expect that, aligned growth is manifested by just piling up of single crystallites along the preferential direction. Whereas, nanorods derived from a precursor of lower pH (=10.8), comprise of several crystallites along a direction normal to the preferential growth, thereby resulting in a larger diameter ( $\sim 19$  nm) of the rods. Considering the lower limit of the nanorod length, the aspect ratio was estimated to be  $\sim 4$  and  $8.5$  corresponding to pH=10.8 and 13.3 cases; respectively (Table 2.3). A higher pH value represents a higher concentration of OH<sup>-</sup> ions and a higher chemical potential of the precursor solution. A higher chemical potential is generally preferable for the growth of higher aspect ratio 1D nanostructures [18]. Here, a nanorod with a higher aspect ratio is obtainable at a higher pH, but along with some amount of inhomogeneity. Figure 2.8(c) shows an enlarged view of a particular region of the nanorod. The clear lattice fringe pattern suggests crystalline nature of the nanorod, giving an interplanar spacing of  $\sim 0.312$  nm and corresponding to the (222) crystallographic plane of the nanosystem. This is in accordance with the XRD pattern of the concerned specimen (pH=10.8). The SAED pattern exhibits a number of arrayed bright spots indicating highly crystalline nature of the specimen under study.

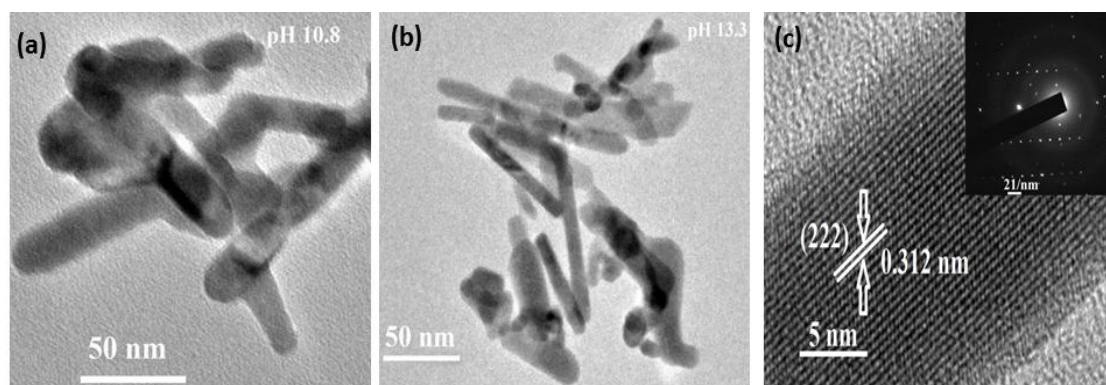


Fig.2.8: (a) and (b) represent HRTEM images of the Gd<sub>2</sub>O<sub>3</sub> nanorods samples synthesized at pH 10.8 and 13.3, respectively. Magnified image of an isolated Gd<sub>2</sub>O<sub>3</sub> nanorod (derived from pH 10.8 precursor) with crystal lattice fringes is depicted in (c). The SAED pattern is being highlighted in the inset of (c).

The growth mechanism of Gd<sub>2</sub>O<sub>3</sub> nanorods has been illustrated schematically, in Figure 2.9. The growth of Gd<sub>2</sub>O<sub>3</sub> nanorods is believed to be governed by a solution-solid phase process [19]. In the hydrothermal process, several drops of aq. NaOH are normally added to adjust the pH of the sol. According to Xu et.al, the growth is neither governed by a catalyst nor directed by any template [20]. On adding aq. NaOH, the Gd<sup>3+</sup> ions of the nitrate product tend to interact with the OH<sup>-</sup> ions in the solution phase. On adding NaOH, a gentle white solution is created which comprised of colloid particles. The colloids, while partially dissolved in water results in a metastable supersaturated solution under hydrothermal environment. On cooling back to room temperature, a large number of seed particles are formed through homogenous nucleation process. It is the seed which becomes the basis of 1D elongated nanostructures formed as a result of coherent re-crystallization process. Accordingly, the whole supersaturated solution is recrystallized to yield highly crystalline rod-shaped structures [20].

With an electronic configuration of [Xe] 4f<sup>7</sup>, Gd<sup>3+</sup> has half-filled shells having a <sup>8</sup>S<sub>7/2</sub> ground state [3]. As we know, the energy absorption in Gd is mediated via transition of 4f electrons to 5d level and then reorganization of the 4f electrons

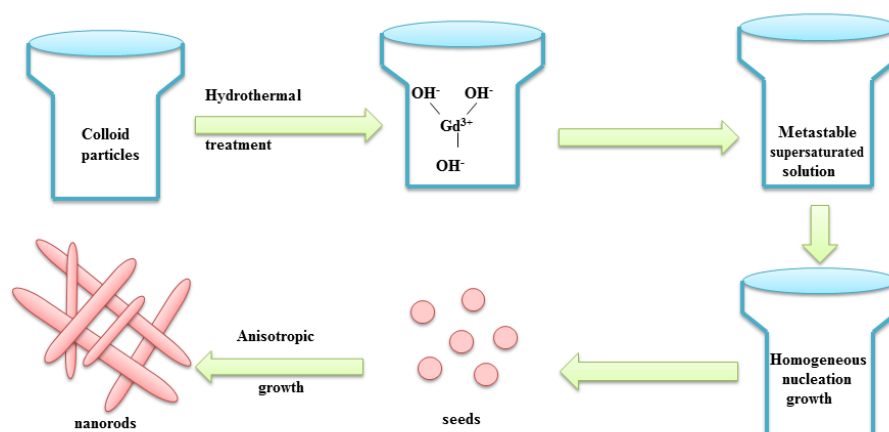


Figure 2.9: A scheme of growth mechanism illustrating formation of  $Gd_2O_3$  nanorods.

into various multiplets. The low lying multiplet above the ground state is  ${}^6P_{7/2}$ . Essentially, two absorption peaks have been observed in the spectra shown in Figure 2.10. The absorption peak positioned at  $\sim 274$  nm is typical for  $Gd_2O_3$  nanorods and can be attributed to  ${}^8S_{7/2} \rightarrow {}^6I_{7/2}$  transitions [3]. The peak positioned at  $\sim 400$  nm, can be assigned to recombination of delocalized electrons close to the conduction band with a single charged state of surface oxygen vacancy as per Wang's proposition [21]. The peak at  $\sim 274$  nm has a strong absorption feature, particularly for the nanorod specimen derived at a higher pH value.

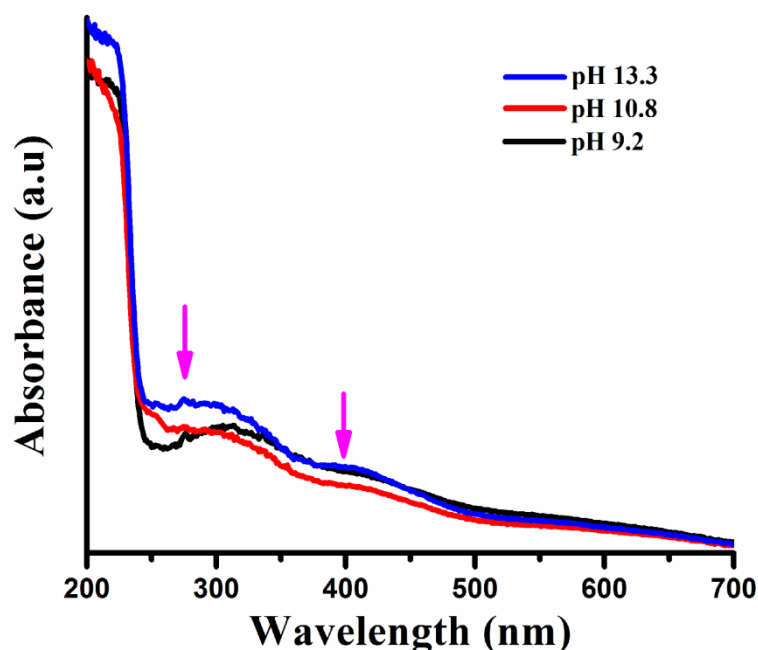


Figure 2.10: Absorption spectra of  $Gd_2O_3$  nanorods synthesized at different pH.

Photoluminescence spectroscopy is a powerful technique to understand the nature of radiative transitions and defect related emission response of a material system. The room temperature emission spectra of the Gd<sub>2</sub>O<sub>3</sub> nanorods, excited at a wavelength of ~ 300 nm are shown in Figures 2.11(a-c). The PL spectra represent characteristic radiative emission features of the nanorods synthesized at different pH values. The asymmetrically broadened spectra signify the superimposition of several peaks of different origin. The deconvolution of each of the spectra through multiple Gaussian fitting could help extracting independent peaks, which can be found in Figures 2.11(a-c). All the samples have showed three prominent peaks, at ~ 413, 442 and 472 nm. The band to band emission was not observable due to the improved defect related emission response. The peak maxima at ~ 413, 442 and 472 nm could be attributed to the surface defects of different Schottky and Frenkel types [11]. The normalized intensity ratio of ~ 413 to 472 nm peak increases with increasing pH value of the desired precursor. In oxide systems, typically oxygen vacancies and interstitials are vastly prevalent and the manifestation of such defects, contribute significantly to the modified luminescence phenomena. Figure 2.11(d) shows a photoluminescence excitation (PLE) spectrum of Gd<sub>2</sub>O<sub>3</sub> nanorods (pH=10.8) monitored at an emission wavelength of ~ 480 nm. PL excitation spectroscopy is a useful method to investigate the electronic structure of the materials with low absorption features. The wavelength of the excitation light is monitored, for a fixed emission wavelength of the material being studied. The peaks, in the PLE spectrum, often represents absorption characteristics of the system under study. Basically, Gd<sup>3+</sup> has a simpler absorption spectrum in comparison to the other trivalent RE ions. Note that, the ground state energy level for Gd<sup>3+</sup> is <sup>8</sup>S<sub>7/2</sub> level. As can be found, the excitation spectra has clearly revealed four important peaks, positioned at ~ 230, 248, 260 and 282 nm. The 230 nm emission peak is attributed to the excitation band of Gd<sub>2</sub>O<sub>3</sub> host [22]. The peak at ~248 nm is ascribed to one of the transitions corresponding to <sup>8</sup>S<sub>7/2</sub> → <sup>6</sup>I<sub>j/2</sub> multiplets, where j=9, 12, 13, 15 and 17 [4]. The



band observed at ~260 nm is assigned to  $^8S_{7/2} \rightarrow ^6D_{9/2}$  transition [23], whereas the peak at ~282 nm corresponded to the  $^8S_{7/2} \rightarrow ^6I_{7/2}$  transition [4].

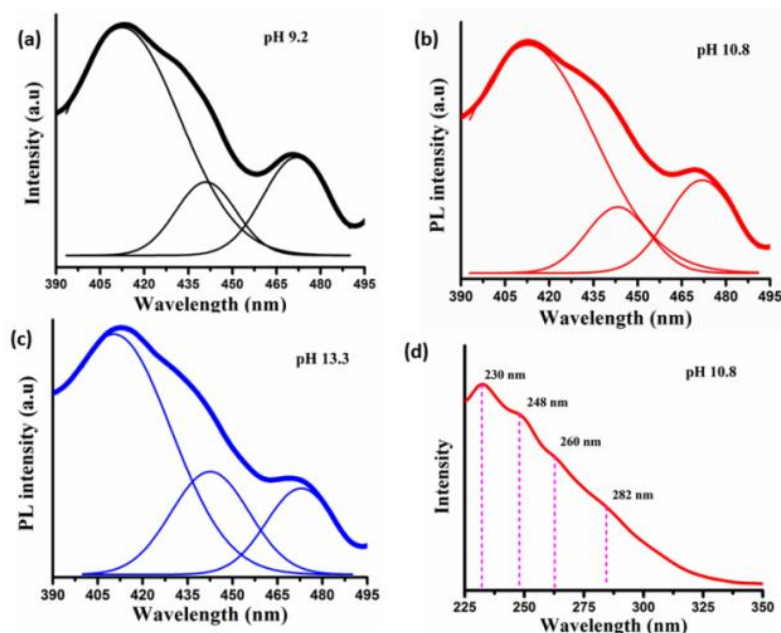


Figure 2.11: Room temperature PL emission spectra  $\lambda_{ex} = 300$  nm with deconvolution of Gd<sub>2</sub>O<sub>3</sub> nanorods obtained from precursors at (a) pH 9.2, (b) pH 10.8 and (c) pH 13.3. a) PL excitation spectrum ( $\lambda_{em} = 480$  nm) for pH=10.8 sample is shown in (d).

Thermogravimetric analysis (TGA) was employed to evaluate the weight loss during progressive heating. Figure 2.12 shows the nature of weight loss in the range of 27-600°C at an atmospheric pressure. In the range 27-250°C, the weight loss (2.11%) is essentially due to the dehydration of absorbed H<sub>2</sub>O in Gd(OH)<sub>3</sub> system. In the range, 250°C-350°C the weight loss is approximately 8.01%, whereas 350°C to 426°C with weight loss 1.3%. The decomposition of Gd(OH)<sub>3</sub> to GdOOH is expected between 250°C -426°C. In final step at 426-594°C with weight loss 2.2% is due to GdOOH to Gd<sub>2</sub>O<sub>3</sub>. The total weight loss measured 13.62%, which almost comparable to the theoretical calculation of Gd(OH)<sub>3</sub> (13.03%) calculated from the reaction of their complete decomposition to produced Gd<sub>2</sub>O<sub>3</sub> [20].

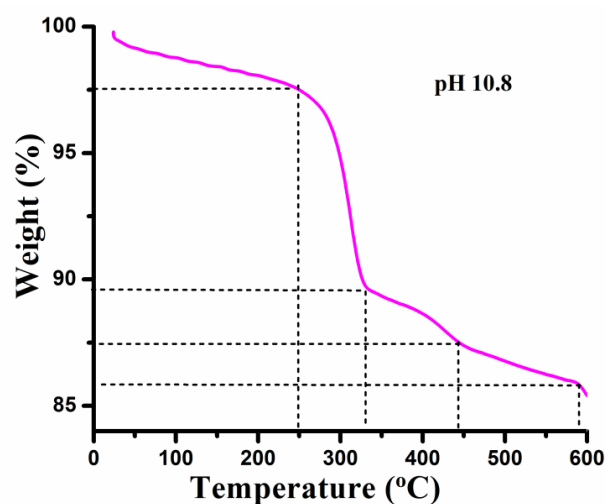


Figure. 2.12: Thermogravimetric analysis of unsintered Gd(OH)<sub>3</sub> product (pH= 10.8).

Raman spectroscopy is a versatile but non-destructive tool to explore vibrational, rotational and other low frequency modes in the systems under study. Raman spectra of the Gd<sub>2</sub>O<sub>3</sub> nanosystems are depicted in Figure 2.13(a). Both the Gd<sub>2</sub>O<sub>3</sub> nanoparticles (GNP) and Gd<sub>2</sub>O<sub>3</sub> nanorods (GNR) exhibited a prominent peak, positioned at  $\sim 360$  cm<sup>-1</sup> and is attributed to the collective signature of  $F_g$  and  $E_g$  modes as predicted in case of the cubic phase Gd<sub>2</sub>O<sub>3</sub> system [24]. While the peak position is retained, the Raman intensity of the GNR system is seen to be improved as compared to the GNP system. Upon Lorentzian fitting of the Raman peak ( $F_g+E_g$  mode), the FWHM corresponding to GNP and GNR systems, are calculated to be  $\sim 25.9$  and  $17.5$  cm<sup>-1</sup>; respectively. The decrease of FWHM of GNR system indicates an enhanced crystallinity along with a stronger phononic response [25]. The phonon life time ( $\tau$ ) of the synthesized samples can also be obtained from the Raman spectra through the energy time uncertainty relation [25]:

$$\frac{1}{\tau} = \frac{\Delta E}{\hbar} = 2\pi c\Gamma, \quad (2.5)$$

where  $\Delta E$  is the uncertainty in the energy of the phonon mode,  $\hbar$  is the reduced Planck's constant and  $\Gamma$  is the FWHM of the Raman peak in unit of cm<sup>-1</sup>. The phonon lifetimes of our GNP and GNR systems are calculated to be 0.20 and 0.32 ps; respectively. Thus, the one-dimensional nanorods (GNRs) are capable

of preserving phononic modes for a longer duration than the modes available in the nanoparticle (GNPs) counterpart.

In order to assess magnetic behaviour of the synthesized product, the Gd<sub>2</sub>O<sub>3</sub> nanorod sample (prepared at a pH of 10.8) was assessed by a Quantum Design MPMS device. Superconducting quantum interference device (SQUID) is a very sensitive magnetometer used to monitor very small changes in magnetic flux and to reveal the magnetic polarization in response to an applied field. The working principle of a SQUID is based on the superconducting loops containing Josephson junctions. The details can be found in literature (Appendix IV) [26]. Figure 8 (b) shows the SQUID magnetometry magnetization vs. magnetic field ( $M-H$ ) plot measured at room temperature (300 K). The magnetization response was measured when applied field was varied in the range -6 T to 6 T. A small hysteresis could be realized initially only upon zooming the central part of the plot. Thus the sample has a dominant paramagnetic phase at room temperature. Even the exact saturation magnetization ( $M_s$ ) could not be predicted. The paramagnetic nature of the nanorods containing Gd<sup>3+</sup> centers is due to the presence of half-filled  $f$  orbitals with seven electrons [4]. Note that, paramagnetic characteristic of a system is caused by a number of magnetic dipole moments summing up to zero in the absence of an applied field. But the net moment becomes nonzero when the system is subjected to a field. However, superparamagnetism arises when thermal energy is strong enough to orient individual moments randomly. As far as superparamagnetism is concerned, in the past, most of the studies were devoted mostly to iron based nanostructures [27]. We anticipate that, the paramagnetism in the Gd<sub>2</sub>O<sub>3</sub> nanorods has arisen from the noninteracting localized nature of the magnetic moments. The magnetic properties of the Gd<sup>3+</sup> ions come from seven unpaired inner  $4f$  electrons, which are closely bound to the nucleus and effectively shielded by the outer closed shell electrons  $5s^25p^6$  from crystal field. Among various paramagnetic RE ions, Gd<sup>3+</sup> ions possess a high magnetic moment due to isotropic electronic ground state  $^8S_{7/2}$ . In order to explore further, we zoomed in the  $M-H$  response near zero field. Remarkably,

the  $M$  exhibits nonzero value in the absence of the field ( $H=0$ ). The coercive field ( $H_c$ ) and remanence ( $M_r$ ) were found to be 0.0066 T and 0.0136 emu/gm; respectively. A low coercive field indicates that, magnetic moments can be reversed even by changing the direction of a weak field.

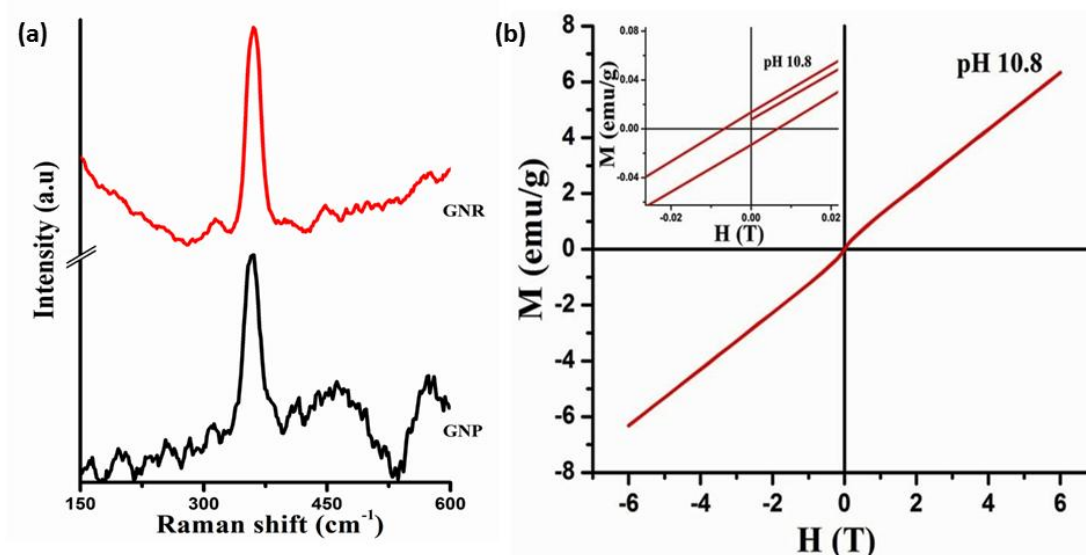


Figure 2.13:(a) Raman spectra of Gd<sub>2</sub>O<sub>3</sub> nanoparticles (GNP) and nanorods (GNR), (b) room temperature  $M \sim H$  plot of the Gd<sub>2</sub>O<sub>3</sub> nanorods synthesized at pH=10.8. The inset shows enlarged view of the plot near zero-field.

### 2.3.2 Synthesis, growth mechanism of Gd<sub>2</sub>O<sub>3</sub> nanorods obtained from nanoparticle seeds

At first, gadolinium acetate hydrate [Gd(CH<sub>3</sub>CO<sub>2</sub>)<sub>3</sub>.H<sub>2</sub>O] (CDH, 99.9% pure) is dispersed in deionized water and then, citric acid [C<sub>6</sub>H<sub>8</sub>O<sub>7</sub>] (CDH 99% pure) is gently added to offer a 1:1 molar solution. The mixture is then subjected to vigorous stirring (~300 rpm) for 30 min, followed by filtration with a whatman paper<sup>®</sup> to yield Gd(OH)<sub>3</sub> product. The Gd(OH)<sub>3</sub> product is then further sintered at a temperature of 800 °C in a tubular furnace, for 4 h, which gave rise to yield Gd<sub>2</sub>O<sub>3</sub> nanoparticles (GNP). The details procedure has been discussed in page no. 37.

Secondly, 0.15 g of the synthesized GNP is dispersed in 35 ml distilled water. The pH is adjusted to ~10 by adding aq. NaOH in a drop wise manner. The mixture is then transferred to a teflon-lined stainless autoclave and

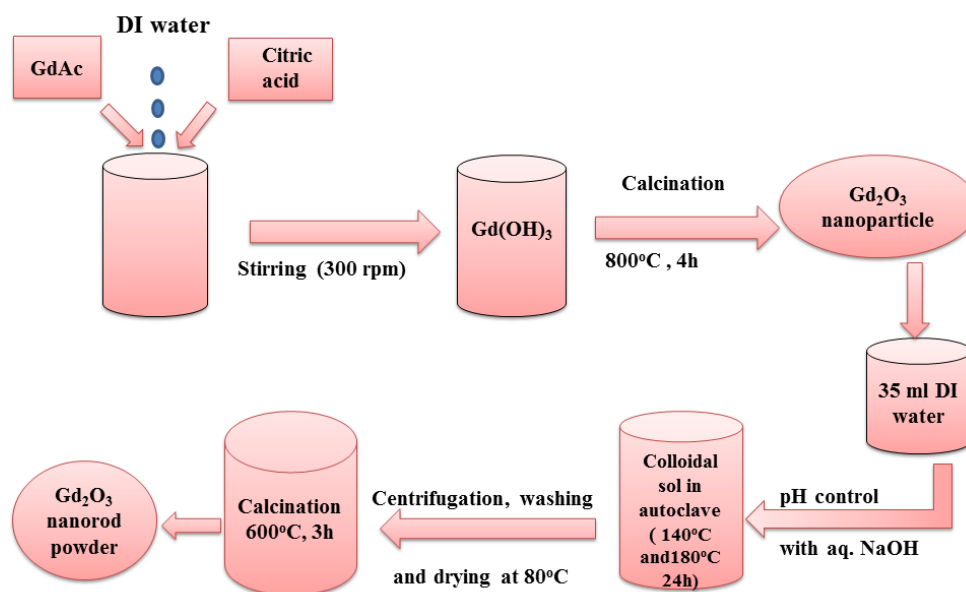


Figure 2.14: Block diagram on various steps involved in the synthesis  $Gd_2O_3$  nanorods

subjected to oven-heating independently while maintaining a temperature of 140 °C and 180 °C, for 24 h. A white solid precipitate of  $Gd_2O_3$  is then collected by filtering the precursor by using a Whatman filter®. The as-received product was washed several times, first with distilled water and then with ethanol, and subsequently followed by oven-drying at a temperature of 80 °C. In order to remove away all the hydroxyl groups present in the system, the powder was finally calcined at a temperature of 600 °C for 3 h. The as- received hydroxyl-free  $Gd_2O_3$  nanorods (GNR) are kept ready for subsequent experimentation. The block diagram of the complete synthesis procedure is shown in Figure 2.14. Figure 2.15 (a) shows a set of XRD patterns of the as-synthesized samples of the  $Gd_2O_3$  nanoparticles and hydrothermally processed  $Gd_2O_3$  nanorods at different temperatures. In each case, the diffraction peaks characterize cubic phase of the  $Gd_2O_3$  system [1]. The distinct peaks corresponded to (222), (400), (440) and (622) crystallographic planes of the samples under investigation ((JCPDS no. 11-0604) [1]. Moreover, we notice a shifting of peaks towards a lower diffraction angle ( $2\theta$ ) in case of  $Gd_2O_3$  nanorods as compared to particles. The shifting of peak position towards a lower diffraction angle would indicate enhancement of lattice parameters [17]. Since (222) appeared as the most

intense peak, most of the crystallites are believed to have preferred orientation along the said plane. Moreover, the average crystallite size ( $d$ ) and microstrain ( $\eta$ ) can be calculated from the W-H expression mentioned in equation (2.1). The GNP system gives an average size ( $d$ ) estimation of  $\sim 11.8$  nm. Whereas, hydrothermally derived GNRs that are processed at a temperature of 140 and 180 °C, and post-calcined at a temperature of 600 °C offered an average crystallite size of  $\sim 12$  and 14 nm; respectively. It may be noted that, without calcination, the GNRs experienced a relatively lower size of  $\sim 8$  nm (Figure 2.15 (a) and Table 2.4). Using the diffraction equation valid for a cubic system (equation no. 2.4) the lattice parameter, ' $a$ ' is estimated to be 1.077, 1.086, 1.091 and 1.086 nm for GNP, GNR without calcination and GNRs (prepared at 140 and 180 °C) subjected to calcination at a higher temperature (Table 2.4). It may be noted that, particle-size dependent lattice expansion is quite common in numerous metal oxides including ABO<sub>3</sub> and ABO<sub>4</sub> types [28, 29]. According to these works, the lattice expansion is due to a proportionately large content of surface defect dipoles which could improve the negative surface stress with the particle size reduction. On the contrary, in our hydrothermally processed Gd<sub>2</sub>O<sub>3</sub> nanosystems, we observed an increased lattice parameter with increasing crystallite size. Our results are consistent to those of a recent article by Yang and co-workers on GdVO<sub>4</sub>:Eu<sup>3+</sup> system [30]. It was proposed that, calcination driven reduced negative surface stress could lead to the desired change [30].

Figure 2.15(b), (c) and (d) depict the HRTEM images of Gd<sub>2</sub>O<sub>3</sub> nanoparticles and nanorods synthesized at a temperature of 140 and 180 °C; respectively. The average size of the nanoparticles, as found from the micrograph, is  $\sim 30$  nm (Figure 2.15(b)). The sample processed at a temperature of 140 °C contained only a few fully grown nanorods along with many underdeveloped rods (Figure 2.15(c)). While exhibiting an approximate diameter of 32 nm, the nanorods as large as 170 nm could be realized. However, as can be evident from Figure 2.15(d), the GNR sample prepared at a temperature of 180 °C, contained a number of fully grown nanorods with occasional hollow

appearances. The average length of the nanorods was found to vary in the range of ~155–190 nm and diameter ~33 nm (Figure 2.15(d)). The inset of Figure 2.15(d) basically depicts a SAED pattern that highlights equi-spaced bright spots arranged in a cubic lattice. Thus we, anticipate mono-crystalline nature of the GNRs synthesized at a temperature of 180 °C.

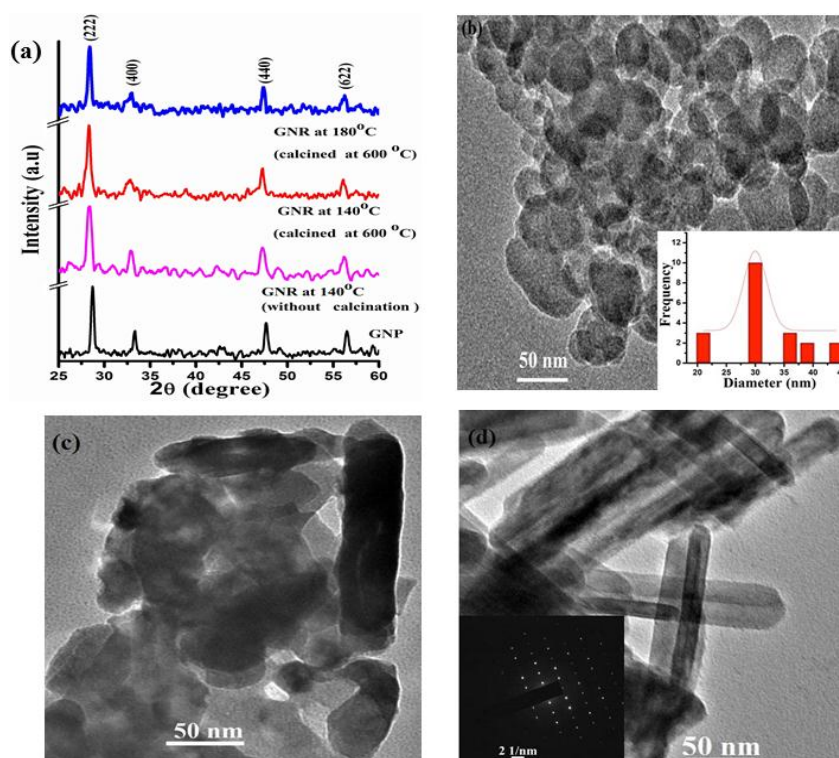


Figure 2.15: (a) XRD patterns of all the Gd<sub>2</sub>O<sub>3</sub> nanosystems. HRTEM images of Gd<sub>2</sub>O<sub>3</sub> (b) seed nanoparticles, and nanorods synthesized at a processing temperature of (c) 140°C, and (d) 180°C. The micrograph shown in (d), characterize well developed nanorods with instances of hollow morphology. The SAED pattern is depicted in the inset of (d).

Table 2.4: Physical characteristics estimated from XRD data

Sample	Average crystallite size ( $d$ ) (nm)	Micro-strain ( $\eta$ ) ( $\times 10^{-3}$ )	Lattice parameter ( $a=b=c$ ) (nm)
GNP	11.8±1.1	-2.2	1.077
GNR 140°C	8±1.2	-5.4	1.086
GNR 140°C (calcined at 600°C)	12±1.4	-3.8	1.091
GNR 180°C (calcined at 600°C)	14±1.6	-1.7	1.086

The presence of fully evolved nanorods at an elevated hydrothermal temperature is attributed to the oriented attachment mechanism [31]. The formation of 1D nanostructure with a higher aspect ratio is essentially possible in presence of higher concentration of OH<sup>-</sup> ions [18]. The reaction temperature is another factor which could affect the aspect ratio of the Gd<sub>2</sub>O<sub>3</sub> nanoscale structures. It is worth mentioning here that, either NaOH or KOH plays an important role in the growth of 1D nanorods. The pH dependent study of Gd<sub>2</sub>O<sub>3</sub> can also be found in the literature [32]. A number of patterns, such as nanospheres, nanorods, nanowires, etc., have been grown through hydrothermal route and with necessary pH control [32,33]. Plate-type morphology are obtained for pH<7, whereas nanorods and nanowires are invariably developed at high pH values. Recently, we have also succeeded in fabricating Gd<sub>2</sub>O<sub>3</sub> nanorods via a catalyst-free hydrothermal approach [34]. The effect of pH has been discussed along with the growth mechanism of the nanorods and apparently, nanorods with high aspect ratios can be revealed in the environment of high pH media. The addition of NaOH and KOH has also been found to help control the phase and morphology of octahedral Fe<sub>3</sub>O<sub>4</sub> crystals as well as  $\alpha$ -Fe<sub>2</sub>O<sub>3</sub> platelets [35, 36].

The growth mechanism of Gd<sub>2</sub>O<sub>3</sub> nanorods is schematically shown in Figure 2.16(a). The corresponding TEM imaging can be found in the lower panel of the same figure. Since the Gd<sub>2</sub>O<sub>3</sub> nanorods are synthesized from the Gd<sub>2</sub>O<sub>3</sub> nanoparticles, we speculate that particle growth would occur along a thermodynamically favourable direction. Previously, melting assisted particle growth was observed due to swift ion impact as well as prolonged sintering [37, 38]. In either case, particle growth has been believed to be accompanied by Ostwald ripening (OR) mechanism. In general, particles obtained by OR mechanism have regular shape and the surface defects produced during growth vanish with the particle dissolution. However, since the model is purely based on the equilibrium condition between the solubilization and precipitation, and less amount of solubility is involved at the nanoscale level, an alternative strategy has been adopted. At nanoscale level, the growth



mechanism, under special circumstance, can be governed by oriented attachment (OA) mechanism [39]. Typically, OA mechanism is based on three factors: (1) dipole-dipole interaction that causes two oriented particles to get attached, (2) interaction prevails until perfect lattice match is ensured, and (3) the adjoining interface eventually vanishes resulting in a single crystal particle [40]. In other words, the nanorod structures can be produced by energetically favourable alignment and coalescence of neighboring particles through elimination of the common boundary. This is clearly evident from the micrographs placed in the middle of the lower panel, shown in Figure 2.16(a). The adjoining particles are indicated by dashed- yellow regions, and the interfaces, indicated by arrow marks in the magnified image.

Since the surface energy plays an important role in the crystal growth process, the growth of elongated nanostructures can occur by lowering the surface energy along a specific direction. The driving force behind anisotropic growth of nanostructures is the crystal surface reactivity [41]. In general, Gd<sub>2</sub>O<sub>3</sub> is a known REO system and possesses a cubic crystal structure. Moreover, it is difficult to grow anisotropic nanostructures through classical OR mechanism. Previously, the OA mechanism was employed to describe the growth process in different oxide systems, such as, TiO<sub>2</sub> [42], CdTe [43], and Y<sub>2</sub>O<sub>3</sub> [41]. In the present situation, at first, the cubic Gd<sub>2</sub>O<sub>3</sub> nanoparticles aggregate due to dipole-dipole interactions, fused gradually and then, recrystallize into single crystalline nanorods. It may be noted that, in our case growth of the nanorods is neither surfactant assisted nor template based but purely governed by dipole-dipole interaction caused by OA.

Under hydrothermal treatment, nanoparticles are randomly oriented following Brownian motion. When two adjacent primary particles collide, the coalescence occurs and that, the particles share a common crystallographic orientation. Thus the two primary particles attach to each other and combine into a secondary one which is still very small and continue to collide and coalesce. The repetition of such an organized event leads to directed assimilation of

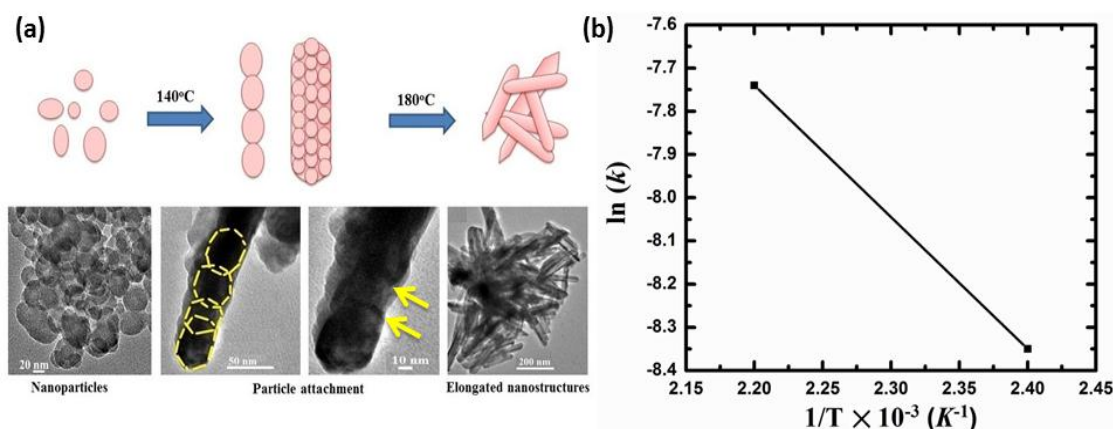


Figure 2.16: (a) Schematic growth mechanism of Gd<sub>2</sub>O<sub>3</sub> nanorods (upper panel) and a set of TEM images illustrating transformation of nanoparticle seeds to nanorods through oriented attachment (lower panel), (b) Arrhenius plot of the reaction constant  $k$  vs. inverse of temperature,  $T^{-1}$ .

particles, resulting in rod shaped structures. The kinetic model governed by OA mechanism can be described by the following reaction [44]:



where  $A_1$  is a primary nanoparticle,  $B$  is the product of the coalescence of two primary nanoparticles and  $k$  is the rate constant involved in the reaction between the two particles. Essentially, reaction occurs between the primary nanoparticles (monomer), whereas nanorods are the collective outcome of the repetitive events. The kinetic model for the OA mechanism can be formulated by [44],

$$d_t = \frac{d_o(\sqrt[3]{2kt+1})}{kt+1} \quad . \quad (2.7)$$

Here,  $d_o$  is the diameter of the precursor nanoparticles,  $d_t$  the diameter of the as grown nanocrystals at time  $t$  and  $k$  is the reaction constant. The values of  $k$  can be calculated for both the reaction temperatures, 140 and 180 °C. In the first case, substituting  $d_t=32$  nm,  $d_o=30$  nm and  $t=1440$  min, in Eq. (2.7),  $k$  is calculated to be  $2.36 \times 10^{-4} \text{ min}^{-1}$ . Using similar kinetics, the rate constant is found to be  $4.34 \times 10^{-4} \text{ min}^{-1}$  for GNRs synthesized at an elevated temperature of 180 °C (diameter,  $d_t \approx 33$  nm). Note that, for evaluating  $k$ , we have not considered the whole length of the nanorods but diameter of the nanorods as

predicted from the TEM analyses. This is because, fully grown elongated structures coexist with the under-developed structures that would represent different stages of growth. The formation of nanorods is believed to be as a result of directed-multi step OA processes dictated by Eq. (2.7). In order to explore the relationship between the kinetic constant,  $k$  and the reaction temperature,  $T$ , the Arrhenius equation can be considered which is given by [44],

$$\ln(k) = -\frac{E_a}{RT} + \ln(A), \quad (2.8)$$

where  $A$  is the pre-exponential factor,  $E_a$  is the activation energy and  $R$  is the universal gas constant (8.314 J·mol<sup>-1</sup> K<sup>-1</sup>). In general, the Arrhenius equation describes the temperature dependence of the rate constant during a chemical reaction. Considering GNP as the base material, one can use the above relation in order to calculate activation energy involved in GNRs, synthesized at different processing temperatures. By plotting calculated  $k$  values in log scale vs.  $1/T$ , we obtained a trace as depicted in Figure 2.16(b). The activation energy involved in the transformation of nanoparticles (GNP) to nanorods (GNR) mediated via oriented attachment in hydrothermal process can be estimated from the slope and is found as, 25.3 kJ/mol. As predicted earlier, (222) crystallographic plane that corresponded to the most dense (111) plane in the cubic cell can become responsible for the preferential growth owing to minimum surface energy conformation [45]. We anticipate that, the nanorod growth, in the present case, is preceded along the said plane following multi OA events in an organized manner (Figure 2.16(a)).

The room temperature photoluminescence spectra of the synthesized samples ( $\lambda_{ex} = 300$  nm) are shown in Figure 2.17(a). The spectra of all the samples are characterized by an intense emission located at ~373 nm. This peak can be attributed to the recombination of the delocalized electron close to the conduction band with a single charged state of the surface oxygen vacancy according to the proposition made by Wang et al. and Zhang et al. [21, 46]. The overall luminescence response of the GNRs is found to be less intense with respect to the GNP counterpart. The deconvoluted PL spectra of the GNP and

the GNR (processed at a temperature of 180 °C) are shown in Fig. 6(b) and (c), respectively. Upon deconvolution, an intense defect emission peak can be found at ~405 nm. After assessment of all the spectra, the intensity ratio between ~373 and 405 nm peaks, is observed to be nearly 1.4 times larger in GNRs than that of GNPs. Not surprisingly, it is the defect mediated emission which is mainly responsible for causing asymmetry in the PL spectra. The right symmetry factor ( $\Delta_R/\Delta$ ), as defined by the ratio of the wavelength share on the right side of the peak maxima to the full width wavelength at half maxima, is slightly improved with the processing temperature. For GNP, GNR at 140°C and GNR at 180°C, the Symmetry factors were calculated to be 0.55, 0.57 and 0.58 respectively. Furthermore, a high concentration of dislocations and point defects can be anticipated when two particles merged into one eliminating the interface in oriented attachment (OA) process [47]. We invoke that, transformation of particles to nanorods may not necessarily offer new kinds of defects but the latter system is likely to possess ample defect centers that account for adequate recombination events.

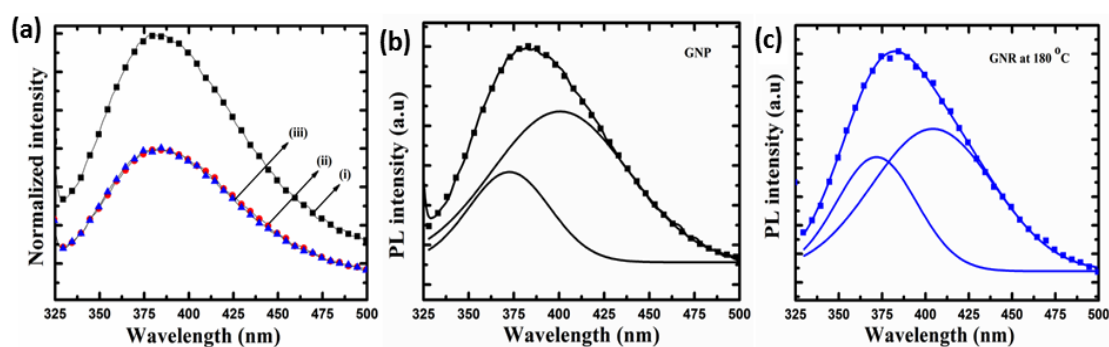


Figure 2.17:(a) PL emission spectra of different Gd<sub>2</sub>O<sub>3</sub> samples: (i) nanoparticles (GNP), and hydrothermally processed nanorods (GNR) at a temperature of (ii) 140 °C and (iii) 180 °C. The deconvoluted PL curves of (b) GNP and (c) GNR systems are shown independently.

Electron paramagnetic resonance (EPR), often referred to as electron spin resonance (ESR), is a widely used spectroscopic technique to study paramagnetic centers in various oxide materials and surfaces [48]. The room temperature, first derivative EPR spectra of the GNP and GNR (processed at

180 °C) systems were recorded in the X-band (9.37 GHz) of the electromagnetic spectrum and with a field modulation of 100 kHz. The variation of signal intensity with the magnetic field is shown in Figure 2.17. The effective  $g$  value was determined from the relation [48]

$$g_{eff} = \frac{hf}{\mu_B H_R}, \quad (2.9)$$

where  $f$  is the frequency ( $\sim 9.37$  GHz) of the microwave source and  $\mu_B$  is the Bohr magneton and  $H_R$  is the resonance magnetic field. The effective  $g$  value is calculated to be  $\sim 2.02$ , which can account for effective hole trapping in the system ( $g > 2$ ) [49]. The respective asymmetry parameters (A/B ratio) for the GNP and GNR systems are  $\sim 1.2$  and 1. The EPR spectra are isotropic in nature. Though vast number of defect types could influence the EPR spectra drastically, the presence of adequate point defects (say, oxygen vacancies) at the edge boundaries keep the spectra symmetric with no further splitting and additional doublet [50].

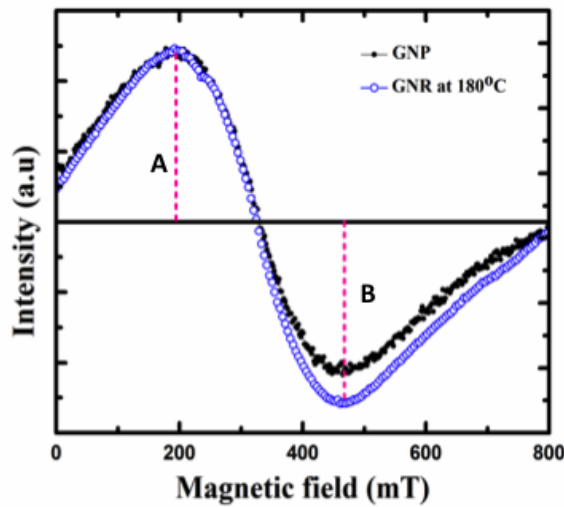


Figure 2.17: Room temperature EPR spectra of the corresponding GNP and GNR systems

## 2.4 Concluding remarks

A variety of nanosystems have been fabricated considering pure and RE doped form of Gd<sub>2</sub>O<sub>3</sub> system. The optical properties of pure and doped Gd<sub>2</sub>O<sub>3</sub> nanosystems have been explored by analyzing UV-Visible absorption and

photoluminescence spectra. The PL spectra of the Tb<sup>3+</sup> doped Gd<sub>2</sub>O<sub>3</sub> nanoparticles were characterized by *D-F* transition along with various native defect mediated emissions. The surfactant free hydrothermal method is found to be suitable for synthesizing of elongated nanostructure of Gd<sub>2</sub>O<sub>3</sub>. XRD studies on Gd<sub>2</sub>O<sub>3</sub> systems have revealed cubic phase of the as-synthesized nanorods while unsintered Gd(OH)<sub>3</sub> product showed a hexagonal phase. HRTEM analyses have revealed a remarkable distribution of nanoscale rods, with aspect ratio nearly doubled when the precursor used was changed from pH = 10.8 to 13.3. Moreover, the evidence of a clear lattice fringe pattern of the rods and ordered bright spots as observable from SAED, suggests that the samples were highly crystalline in nature. TGA analysis on the hydroxide product has revealed the total weight loss ~ 13.62% up to a temperature of ~ 600 °C. The prominent peak in the Raman spectra attributed to the cubic phase of the nanorods. The magnetic measurements have revealed paramagnetic nature of the Gd<sub>2</sub>O<sub>3</sub> nanorods at room temperature. On the other hand, the cubic phase, Gd<sub>2</sub>O<sub>3</sub> nanorods were synthesized from nanoparticles via a surfactant free hydrothermal method. Mediated by oriented attachment (OA) mechanism, a complete growth of nanorods was realized when the hydrothermal reaction was carried out at a temperature of 180 °C. The nanorods were expected to have formed through multi-step OA, accompanied with an activation energy of 25 kJ/mol. The EPR spectrum of the nanorods was more isotropic than that of the nanoparticles, while both gave  $g_{eff} > 2$ .

## References

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