

## **Chapter 4.**

**Power dependent Raman spectra, nitrogen adsorption-desorption isotherms, and photocatalytic activity of WS<sub>2</sub> nanosystems exposed to UV- and  $\beta$ -radiation**



## 4.1 Introduction

The organic and inorganic pollutants found in industrial wastewater have posed significant health risks, especially in developing countries, in recent decades. The textile and paper industries use vast quantities of organic dyes and synthetic pigments, leading to severe environmental issues. Synthetic dyes are commonly employed in industries such as textiles, leather, cosmetics, and pulp and paper, resulting in the disposal of alarming amounts of coloured organic dyes. These dye residues are highly carcinogenic, and a significant portion (1% -10%) of the dye is wasted during the product dyeing process, eventually ending up in wastewater. Specifically, azo dyes are the primary constituents of wastewater from textile industries and processing units [1,2].

Two-dimensional (2D) transition metal dichalcogenides (TMDCs) have gained considerable attention as photocatalysts due to their strong light absorption, remarkable chemical stability, and cost-effectiveness. These materials have been widely employed in pollutant removal [3,4]. It is well established that the edge sites of TMDCs play a crucial role in catalytic applications, as they provide active surface sites, while the basal plane remains relatively inert [5]. Additionally, radiation-based modifications using radionuclides have emerged as a valuable and cost-effective approach for enhancing the properties of TMDC materials [6]. Specifically, exposure to charged particle ions, energetic electrons, and  $\gamma$ -rays induces chalcogen vacancy-related point defects within TMDCs, inducing observable modification to the material properties at large. Furthermore, the microporous nature of 2D TMDC materials has demonstrated immense potential in hydrogen evolution reaction, photocatalysis, and electrocatalysis, along with nonlinear light absorption capabilities [7]. One of the effective strategies for improving nanomaterial performance is exploiting porous 2D layered structures. These hierarchical 2D materials exhibit superior physical and chemical properties due to their distinctive structural and morphological advantages. These porous hierarchical materials exhibit multiple ranges of porosity across different length scales, including (i) micropores (dia.  $<2$  nm), (ii) mesopores (dia.  $\sim 2$ -50 nm), and (iii) macropores (dia.  $>50$  nm) [8]. These materials also exhibit bimodal pore distributions, such as (i) micro-meso, (ii) meso-macro, and (iii) micro-macro, as well as trimodal distributions encompassing micro-to-macro and meso-to-macro pores. The ability to control pore size distribution is particularly significant for applications requiring specific functionalities [9,10]. According to reports, microporous



structures effectively harness visible light due to their increased surface area, which enhances photo-induced electron transfer. This, in line, leads to improved photocatalytic performance [11,12].

This chapter examines the impact of radiation exposure on layered WS<sub>2</sub> systems using (UV) and beta ( $\beta$ ) rays. Irradiation on WS<sub>2</sub> systems is employed to develop a porous structure with diverse micro- and mesopores. Furthermore, the effect of these porous structures on the photocatalytic dye degradation activity of congo red, an azo dye, was investigated using irradiated WS<sub>2</sub> systems as a nano-catalyst.

#### 4.2 Irradiation of WS<sub>2</sub> system with UV- and $\beta$ -rays

The exfoliated WS<sub>2</sub> powder was subjected to irradiation utilising UV- and  $\beta$ -rays. For this process, 1.5 g of exfoliated WS<sub>2</sub> material was evenly spread onto microscope glass slides for each experimental condition. One of these slides was placed inside a UV chamber, positioned 12 cm below a UV lamp offering excitation wavelength,  $\lambda_{\text{ex}} = 365$  nm, with an intensity of  $\sim 800 \mu\text{W}/\text{cm}^2$  and exposed for a time duration of  $\sim 30$  min. Similarly, another 1.5 g sample of exfoliated WS<sub>2</sub> powder was irradiated with  $\beta$ -rays inside a Pb-shielded chamber. A laboratory-grade <sup>137</sup>Cs source was employed to emit  $\beta$ -radiation, transitioning through a metastable energy state of  $\sim 0.51$  MeV and a stable state near  $\sim 1.17$  MeV, delivering a total radiation dose of 0.2 kGy (shown in *Appendix* (Fig. A2)) [13]. All irradiation procedures were conducted at room temperature (300 K).

#### 4.3 Execution of photocatalytic experiment

The photocatalytic activity was evaluated using commercially available congo red (CR) dye (ACS reagent, Himedia<sup>®</sup>) as the target material. Photocatalytic experiments were conducted on three different samples, specifically exfoliated WS<sub>2</sub>, UV-irradiated WS<sub>2</sub>, and  $\beta$ -irradiated WS<sub>2</sub>. To prepare the stock solution, 25 mg/L of CR dye was mixed in a borosil beaker, and 125 mg/L of exfoliated WS<sub>2</sub> was added as a photocatalyst. The mixture was stirred using a magnetic stirrer at room temperature in complete darkness for  $\sim 1$  h to achieve adsorption-desorption equilibrium. Similarly, CR dye solutions with UV- and  $\beta$ -irradiated WS<sub>2</sub> catalysts were prepared to assess their photocatalytic degradation efficiency. The prepared solutions were placed inside a UV chamber, with the light source being positioned 12 cm above the specimen under study. The exfoliated WS<sub>2</sub> sample was exposed for approximately 75 min, while the irradiated WS<sub>2</sub> samples were exposed for a



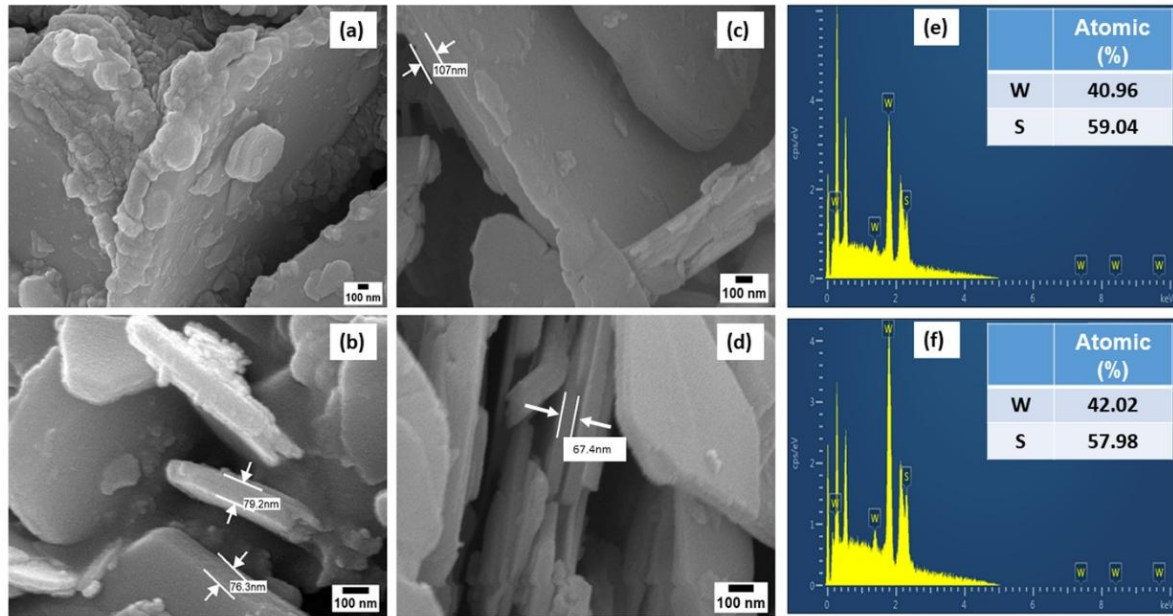
time duration of  $\sim 45$  min. At an interval of 15 min, 3 mL of the solution was collected in a quartz cuvette to measure the absorbance spectra. The degradation percentage of the dye was calculated using the popular formula [14],

$$\% \text{ degradation} = \left(1 - \frac{C}{C_0}\right) \times 100 = \left(1 - \frac{A}{A_0}\right) \times 100 \quad (4.1)$$

where  $C_0$  and  $C$  represent the concentration of CR dye at an initial time (0 min) and time  $t$ , respectively, while  $A_0$  and  $A$  correspond to the strength of the maximum peak in the recorded absorbance spectra.

#### 4.4. Morphological and elemental analysis

The surface morphologies of bulk WS<sub>2</sub>, exfoliated WS<sub>2</sub>, and systems subjected to UV- and  $\beta$ -irradiation are analysed with the help of captured FE-SEM images, depicted in Fig. 4.1(a-d). The bulk WS<sub>2</sub> appears as thick, agglomerated stacks of layers. In contrast, exfoliated WS<sub>2</sub> and irradiated systems exhibit a distinct separation of thin sheets while preserving



**Figure 4.1:** FE-SEM images of (a) bulk WS<sub>2</sub>, (b) exfoliated WS<sub>2</sub>, (c) UV-irradiated WS<sub>2</sub> and (d)  $\beta$ -irradiated WS<sub>2</sub> at lower magnifications (scale bar: 100 nm) respectively can be seen, while higher magnifications of the respective systems are shown in the right corner of the figures as insets; (e, f) EDX spectra of UV- and  $\beta$ -irradiated WS<sub>2</sub> systems.



their layered nature. Notably, some fragmented or uneven segments may have formed due to the shear forces acting on the sheets during prolonged irradiation. Elemental analysis of the irradiated systems, as shown in Fig. 4.1(e, f), revealed sulfur vacancies caused by exposure to UV and  $\beta$ -rays. By calculating the sulfur to tungsten (S/W) ratio from EDX spectra of scanned areas, a further decline in the stoichiometric ratio of WS<sub>2</sub> was observed, reaching a W:S ratio of 1:1.44 for UV exposure and 1:1.38 for  $\beta$ -ray exposure cases. The decrease in the number of sulfur atoms in the WS<sub>2</sub> system after  $\beta$ -irradiation is attributed to the much higher energy of beta particles, which can remove a larger number of sulfur atoms than the lower-energy UV irradiation.

#### 4.5. Laser power-dependent Raman analysis

Raman spectra were obtained for bulk WS<sub>2</sub>, exfoliated WS<sub>2</sub>, and UV- and  $\beta$ -irradiated WS<sub>2</sub> systems, as shown in Fig. 4.2(a). The two most prominent peaks correspond to the vibrational mode observed at  $\sim 351 \text{ cm}^{-1}$ , which emerges from the strong covalent bonding between S and W atoms, while the out-of-plane mode, at  $\sim 421 \text{ cm}^{-1}$ , corresponds to vibrations amongst layers of S atoms linked through van der Waals bonding. Additionally, poorly resolved second-order or mixed modes were located at  $\sim 523 \text{ cm}^{-1}$  and  $\sim 583 \text{ cm}^{-1}$ , which are attributed to the  $E'_{2g} + LA$  and  $A_{1g} + LA$  modes, respectively [15]. Deconvoluted spectra of the WS<sub>2</sub> system after  $\beta$ -irradiation is shown in the *Appendix* (Fig. A3). A noticeable redshift and broadening of these peaks, associated with both in-plane and out-of-plane vibrational modes, suggest a reduction in the number of layers after exfoliation compared to the pristine structure. This can be realized from an increase in the intensity ratio of  $E'_{2g}$ -to- $A_{1g}$  Raman modes after exfoliation, followed by UV treatment. However, in the case of exfoliated WS<sub>2</sub> exposed to high-energy  $\beta$ -rays, an anomalous behaviour is observed due to the presence of adequate sulfur vacancies, leading to the disruption of W-S bonds (Fig. 4.2(b)). Additionally, the phonon lifetime ( $\tau$ ) was determined using the uncertainty principle from the acquired Raman spectra, following the equation [16],

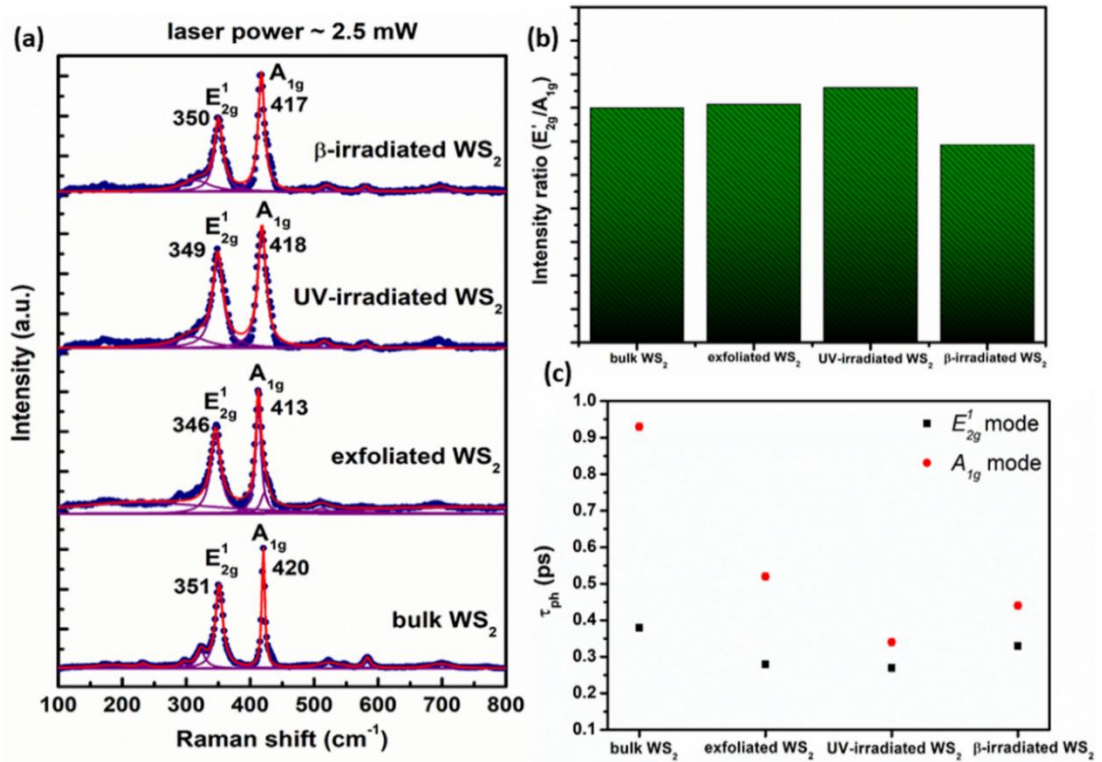
$$\frac{1}{\tau} = \frac{\Delta E}{\hbar} = 2\pi c\Gamma \quad (4.2)$$

where  $\Delta E$  represents the phonon mode energy,  $c$  is the speed of light ( $3 \times 10^8 \text{ m/s}$ ),  $\Gamma$  denotes the full width at half maximum (FWHM) of the Raman peak (obtained through



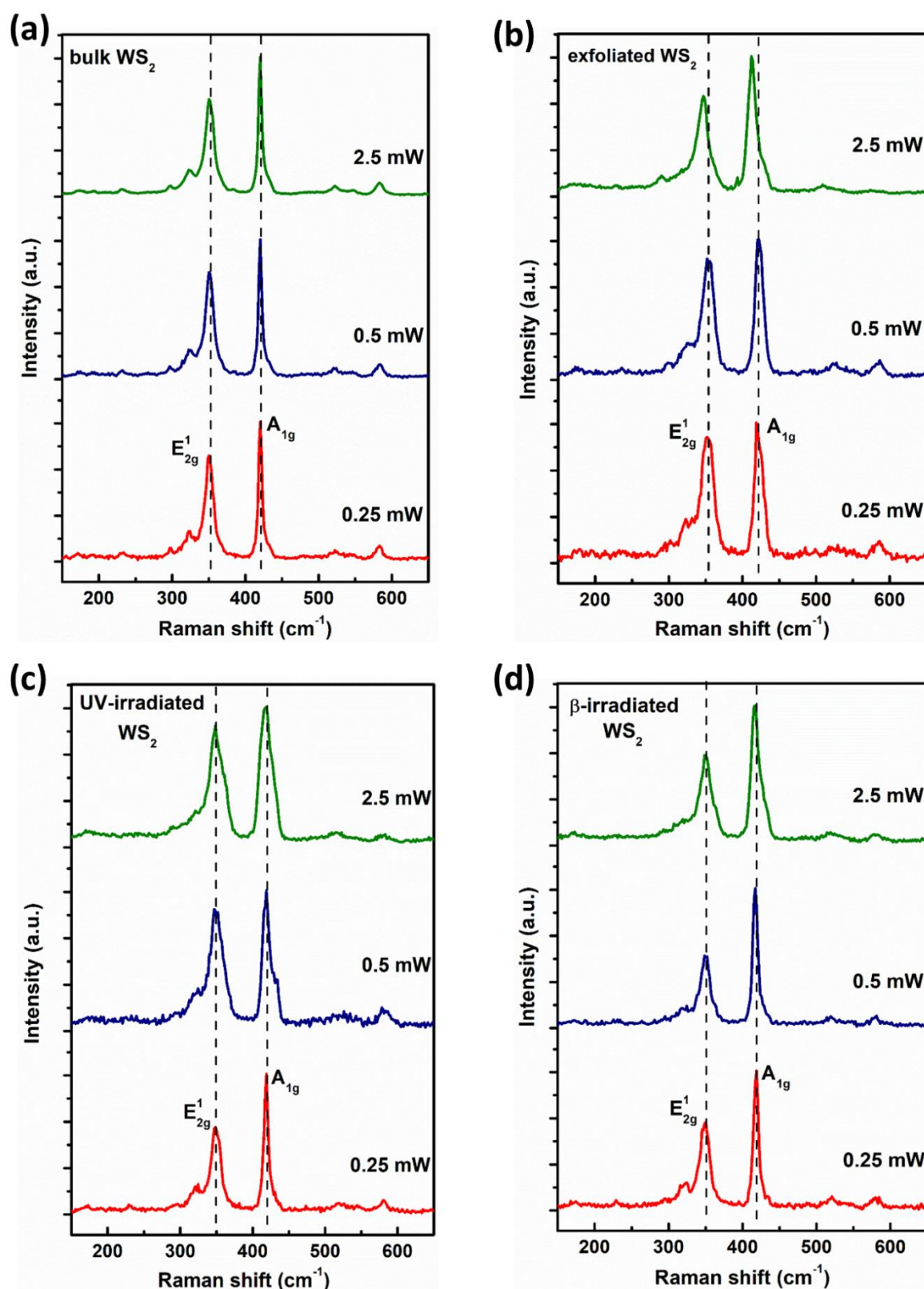
multi-peak Lorentzian fitting of the  $E_{2g}^1$  and  $A_{1g}$  modes), and  $\hbar$  is reduced Planck's constant. A comparison of the characteristic phonon lifetimes for pristine and irradiated systems calculated using equation 4.2 is presented in Fig. 4.2(c). The peak positions of the prominent  $E_{2g}^1$  and  $A_{1g}$  modes, along with their intensity ratio ( $E_{2g}^1/A_{1g}$ ), FWHM, and phonon lifetime ( $\tau$ ), are summarized in Table 4.1.

Raman spectra were also plotted for bulk WS<sub>2</sub>, exfoliated WS<sub>2</sub>, and UV- and  $\beta$ -irradiated WS<sub>2</sub> nanosystems, with laser power excitation ranging from 0.25 to 2.5 mW (Fig. 4.3). While variations in vibrational intensities were observed with different laser power excitations, there was no significant peak shift in the studied systems, except for the exfoliated WS<sub>2</sub>, which exhibited a redshift in prominent Raman peaks at higher laser excitation ( $\sim 2.5$  mW). Furthermore, mixed or second-order modes, explicitly the  $E_{2g}^1 + LA$  and  $A_{1g} + LA$  modes, which appear as weakly resolved peaks at  $\sim 523$  cm<sup>-1</sup> and  $\sim 583$  cm<sup>-1</sup>, respectively, were suppressed with increasing laser power in the exfoliated sample as well



**Figure 4.2:** (a) Raman spectra, (b) histograms representing the  $E_{2g}^1$  to  $A_{1g}$  intensity ratio of bulk WS<sub>2</sub>, exfoliated WS<sub>2</sub>, UV- and  $\beta$ -irradiated WS<sub>2</sub> systems. (c) Plot showing the variations of phonon lifetimes vs. vibrational modes of the systems.



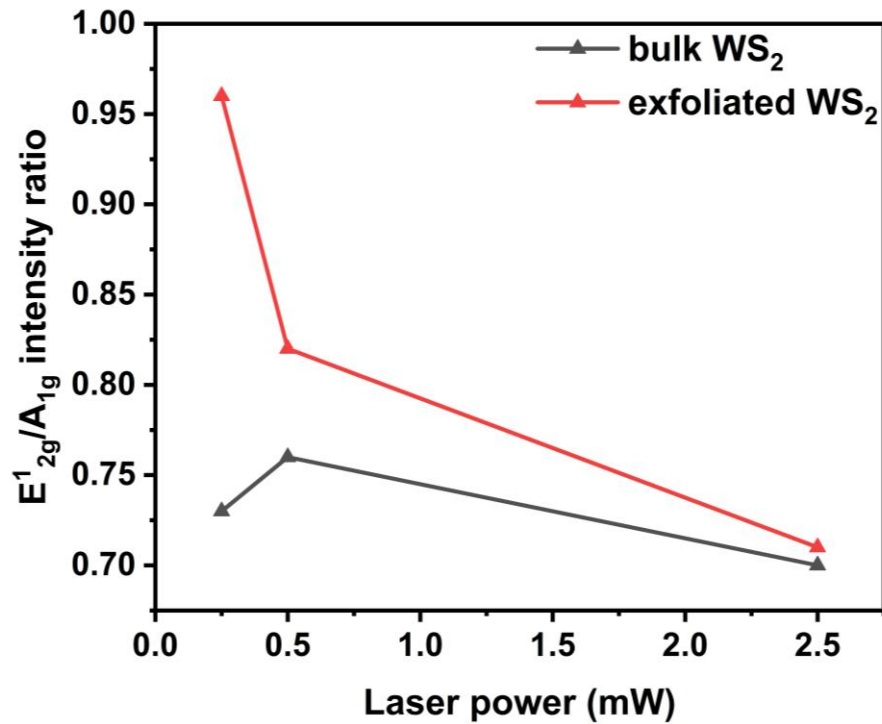


**Figure 4.3:** Raman spectra of (a) bulk WS<sub>2</sub>, (b) exfoliated WS<sub>2</sub>, (c) UV- and (d)  $\beta$ -irradiated exfoliated WS<sub>2</sub> nanosystems acquired for different laser powers ( $\sim$ 0.25-2.5 mW).



**Table 4.1.** Raman mode assignment and parameters obtained from the Raman spectra acquired for the laser excitation power of ~2.5 mW.

Sl. No.	Sample	Peak position (cm <sup>-1</sup> )		$E^I_{2g}/A_{1g}$ (intensity ratio)	FWHM modes (cm <sup>-1</sup> )		Phonon lifetime, $\tau$ (ps)	
		$E^I_{2g}$ mode	$A_{1g}$ mode		$E^I_{2g}$	$A_{1g}$	$E^I_{2g}$ mode	$A_{1g}$ mode
1	Bulk WS <sub>2</sub>	351	421	0.70	13.8	5.7	0.4	0.9
2	Exfoliated WS <sub>2</sub>	346	413	0.71	18.8	10.2	0.3	0.5
3	UV-irradiated WS <sub>2</sub>	349	418	0.76	19.9	15.4	0.3	0.3
4	$\beta$ -irradiated WS <sub>2</sub>	350	417	0.59	16.3	12.0	0.3	0.4

**Figure 4.4:** Laser power excitation vs.  $E^I_{2g}$ -to- $A_{1g}$  intensity ratio for (a) bulk WS<sub>2</sub> and (b) exfoliated WS<sub>2</sub> nanosheets.

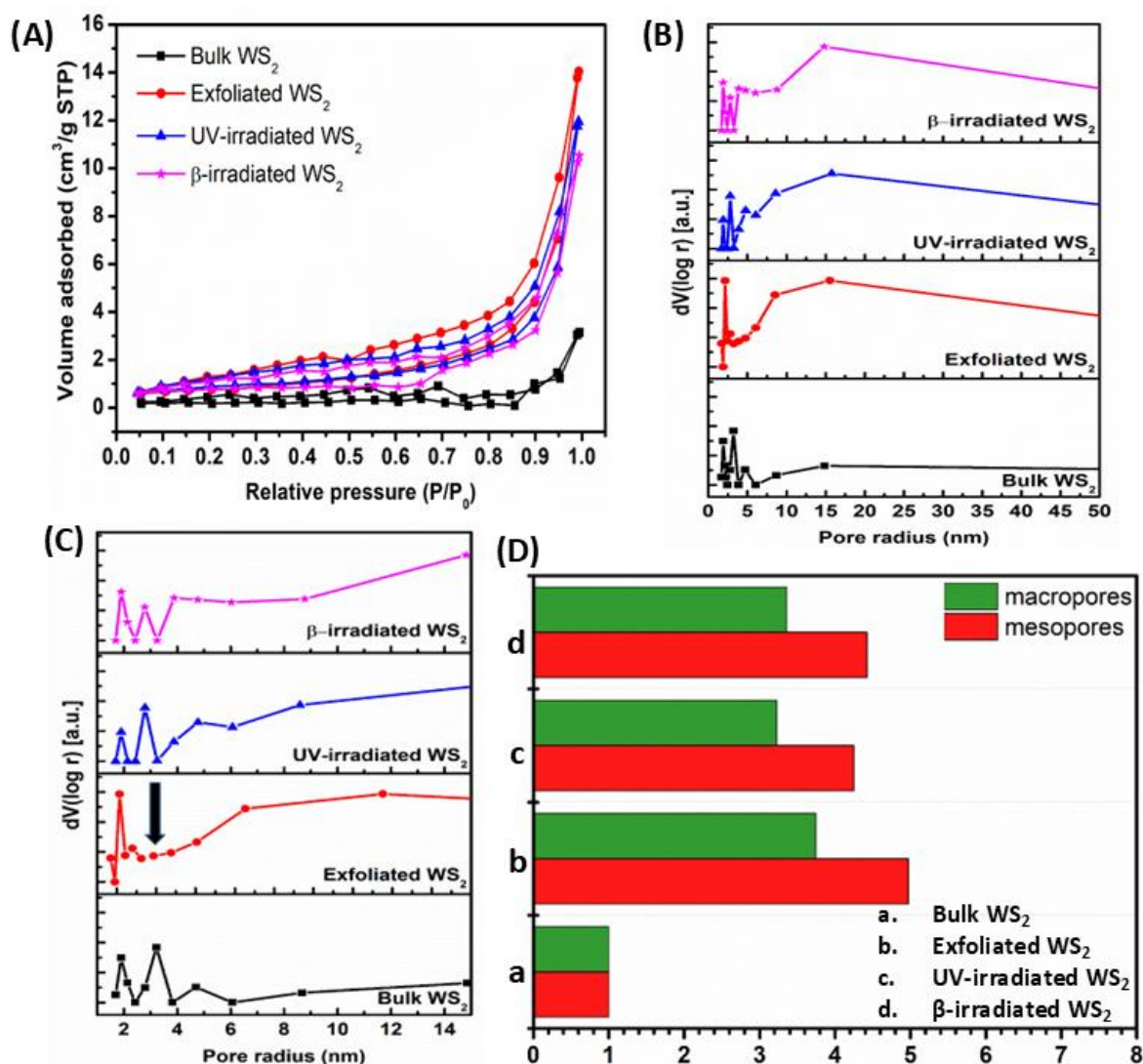


as after UV- and  $\beta$ -irradiation. Additionally, the laser power-dependent Raman observations ( $\sim 0.25$ – $2.5$  mW) concerning the intensity ratio of the  $E'_{2g}$ -to- $A_{1g}$  modes for both bulk and exfoliated WS<sub>2</sub> nanosheets can be found in Fig. 4.4. Notably, the intensity ratio in exfoliated WS<sub>2</sub> is significantly higher than in bulk WS<sub>2</sub>, particularly at lower laser power ( $\sim 0.25$  mW). The variations in Raman intensity observed in exfoliated WS<sub>2</sub> with changing laser power can also be attributed to optical interference effects, which are more prominent in thinner layers. In contrast, for bulk WS<sub>2</sub> with a greater number of stacked layers, the Raman laser light has a negligible optical interference effect due to the thickness of the material [17]. However, as the laser power increases, the intensity ratio gradually decreases, aligning with the vibrational characteristics of the bulk system. This behaviour may be attributed to localized heating effects at higher laser powers, which would weaken the atomic bonds [18]. A redshift in the  $E'_{2g}$  mode of the exfoliated system, prominent for laser power excitation of 2.5 mW, as depicted in Fig. 4.3(b), suggests the existence of a microstrain in the system. It is possible that this strain contributes to the noted drop in intensity ratio with increasing laser power [19].

#### 4.6. Nitrogen adsorption/desorption process and pore-size analysis

The adsorption and desorption isotherms measured at liquid nitrogen (N<sub>2</sub>) temperature ( $\sim 77$  K) have been analysed to determine the pore-size distribution in bulk WS<sub>2</sub>, exfoliated WS<sub>2</sub> and WS<sub>2</sub> system exposed to UV light at an excitation wavelength of  $\sim 365$  nm for  $\sim 30$  min and  $\beta$ -rays (delivering a dose of 0.2 kGy). The N<sub>2</sub> adsorption/desorption isotherms and the corresponding pore-size distributions for these systems are presented in Fig. 4.5. The adsorption isotherm follows the principles of gas adsorption, where N<sub>2</sub> molecules are introduced into the system as pressure increases. In contrast, desorption occurs as the pressure is released, often resulting in a hysteresis loop. To be mentioned, pore diameters larger than  $\sim 4$  nm usually characterize a hysteretic trend [20]. A significant increase in the volume of adsorbed gas at higher relative pressures ( $p/p_0=0.5$ - $1.0$ ) produces a type-IV isotherm, indicating the presence of a large number of mesopores (dia. 2-50 nm), as shown in Figure 4.5(a). This behaviour is attributed to capillary condensation within these mesopores [21]. As well, the isotherms display a type-H3 hysteresis loop, as classified by IUPAC, suggesting the presence of slit-shaped pores [22] and indicating the absence of limiting adsorption at higher  $p/p_0$  values [23]. Furthermore, as seen in Fig. 4.5(A), the linear segments of the exfoliated and irradiated samples have steeper slopes compared to





**Figure 4.5:** (A) N<sub>2</sub> adsorption-desorption isotherms, (B) pore size distribution plots, and (C) zoomed view of pore size distributions of bulk WS<sub>2</sub>, exfoliated WS<sub>2</sub>, UV- and  $\beta$ -irradiated WS<sub>2</sub>, respectively. (D) Comparative bars on meso- and macro-pore size distribution of the specimens considered by comparing the area under the respective curves in the range,  $d \sim 8-100$  nm.

bulk WS<sub>2</sub>, suggesting a rapid yet uniform adsorption/desorption process in the exfoliated, UV-irradiated, and  $\beta$ -irradiated WS<sub>2</sub> systems.

The pore diameters and volume of the porous structures were analysed using the Barrett–Joyner–Halenda (BJH) model [24]. The BJH pore-size distribution graphs for bulk and exfoliated WS<sub>2</sub>, and WS<sub>2</sub> systems exposed to UV and  $\beta$ -irradiation are shown in Fig. 4.5 (B,C). The three distinct peaks at  $\sim 3.8$  nm, 6.4 nm, and 9.4 nm indicate the coexistence of various mesopores in the bulk material, apart from a minimal presence of macropores.



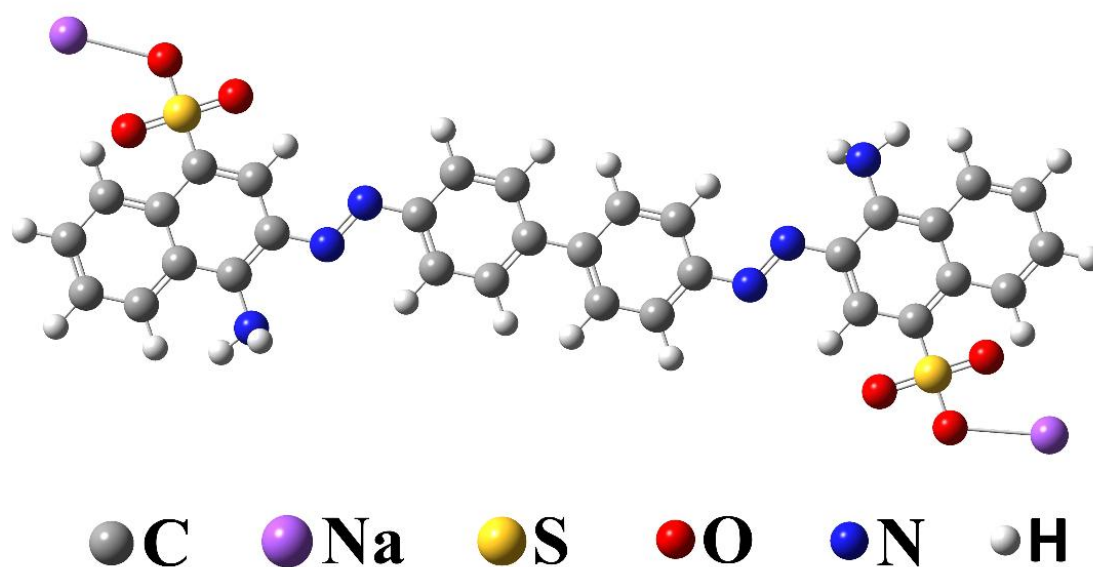
However, in exfoliated WS<sub>2</sub>, the absence of the  $\sim 6.4$  nm peak suggests the formation of new pore sizes, likely due to the shrinking of individual pores or the merging of mesopores into larger ones. This results in a wider range of mesopore size distribution. Interestingly, samples subjected to UV- and  $\beta$ -irradiation exhibit a wider range of pore sizes, spanning from mesopores to macropores, along with the recurrence of the mesopores size of  $\sim 6.4$  nm. In these irradiated systems, ample mesopores are developed and exhibit greater size variation, while macropores are also present in adequate amounts. This is evident from the area under the curves considered in the range of 8-100 nm (shown in *Appendix* (Fig. A4)), which includes both mesopores (8-50 nm) and macropores (50-100 nm), as illustrated in Fig. 4.5 (D). Notably, the analysis of macropores was limited to pore sizes up to 100 nm due to the lack of sufficient data points beyond this range.

#### 4.7 Congo red dye degradation activity

The photocatalytic degradation efficiency of exfoliated WS<sub>2</sub> and a comparative analysis of exfoliated and irradiated WS<sub>2</sub> systems on congo red (CR) dye under UV light exposure were assessed using UV-vis spectral features. As is known to most chemists and researchers, CR is a water-soluble azo dye that appears as a red colloidal solution, composed of a phenyl ring with two identical halves connected to a naphthalene moiety. It is the sodium salt of benzidine-diazo-bis-1-naphthylamine-4-sulfonic acid, with the chemical formula C<sub>32</sub>H<sub>22</sub>N<sub>6</sub>Na<sub>2</sub>O<sub>6</sub>S<sub>2</sub>. The molecular structure of CR dye is presented in Fig. 4.6. This azo dye, which belongs to the largest class of organic dyestuffs, contains a chromophoric  $-N=N-$  azo group [25]. The CR dye exhibits a strong absorbance peak at  $\sim 497$  nm with a broad spectral feature, primarily absorbing light in the blue region. Under UV light exposure, the dye undergoes rapid degradation through cleavage of the  $-N=N-$  bond when exfoliated WS<sub>2</sub> is introduced as a catalyst, exposed up to  $\sim 75$  min (Fig. 4.7(a)). Additionally, a comparative analysis was conducted for exfoliated WS<sub>2</sub> and UV- and  $\beta$ -irradiated WS<sub>2</sub> catalysts on CR dye over durations of  $\sim 15$  min and  $\sim 45$  min, respectively (Fig. 4.7(b, c)). While the exfoliated WS<sub>2</sub> and UV-irradiated WS<sub>2</sub> system displayed inadequate photocatalytic degradation efficiency of only 8.2% and 10.2 % respectively, till 45 min of UV light exposure. Similarly, WS<sub>2</sub> nanostructures showed poor performance, degrading just 15% of the harmful dye even after 105 min [26]. On the other hand, the WS<sub>2</sub> system irradiated with  $\beta$ -rays demonstrated enhanced photocatalytic degradation of CR dye, achieving degradation efficiencies of around 58.6 % and 59.4 % under UV

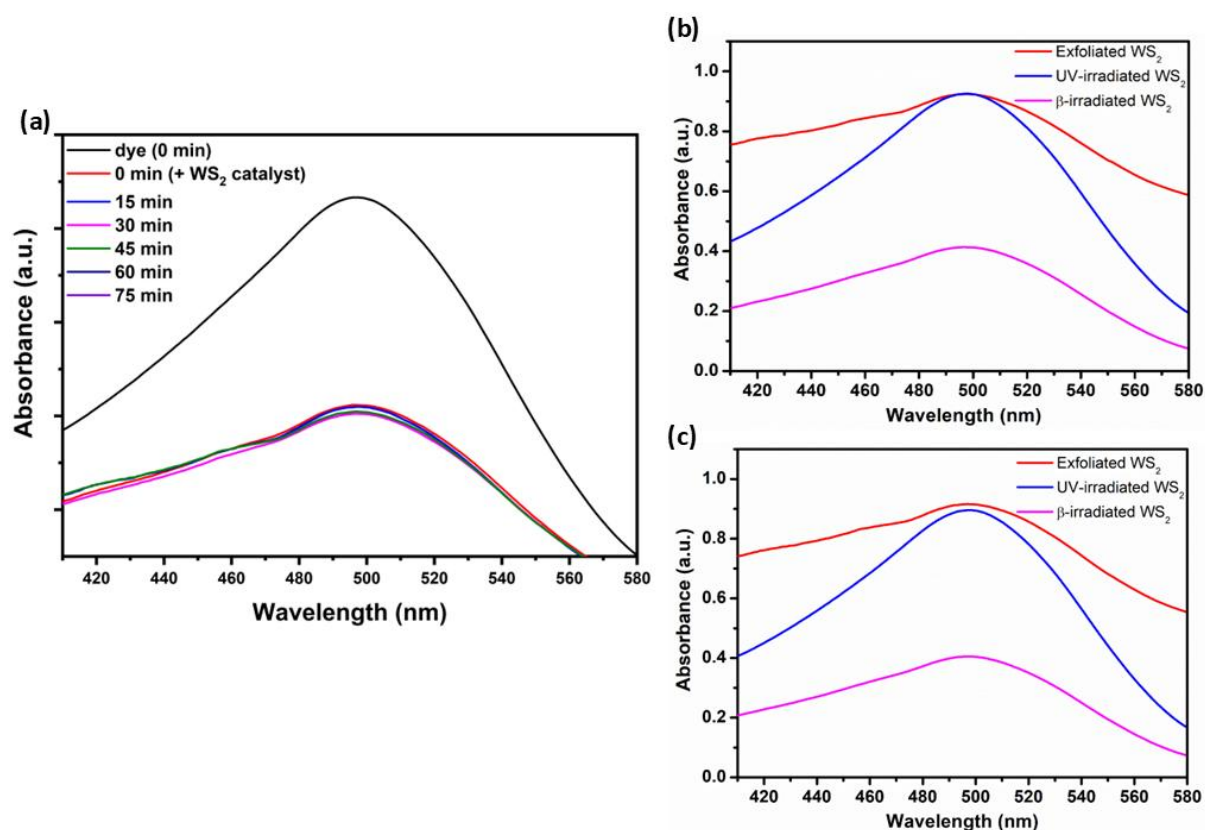


illumination for  $\sim 15$  and  $\sim 45$  min, respectively, as summarized in Table 4.2. This improved performance can be attributed to the presence of varied pores and a relatively higher concentration of sulfur atomic vacancies in the  $\beta$ -irradiated WS<sub>2</sub> system, as revealed by atomic percentage calculations from the EDX spectra. The  $\beta$ -irradiation, a form of ionizing radiation, can deliberately introduce a high number of defects into WS<sub>2</sub> nanocatalyst due to its relatively high energy, significantly greater than that of UV irradiation, which acts in the eV range. These defects can alter the electronic structure of the catalyst or create new active sites, thereby enhancing its catalytic performance [27]. Such structural modifications in the WS<sub>2</sub> catalyst promote the efficient separation and transfer of photo-generated charge carriers while simultaneously reducing the recombination rate of electron-hole pairs. Moreover, sulfur defects are known to enhance the activation of O<sub>2</sub>, leading to the generation of more reactive O<sub>2</sub><sup>-</sup> and  $\cdot$ OH species [28]. As a result, the overall photocatalytic efficiency of the material is significantly improved. The possible mechanism of the photocatalytic degradation process of CR dye with WS<sub>2</sub> material as a photocatalyst under UV light illumination is presented in Fig. 4.8.



**Figure 4.6:** The molecular structure of congo red dye in stick-ball configuration.



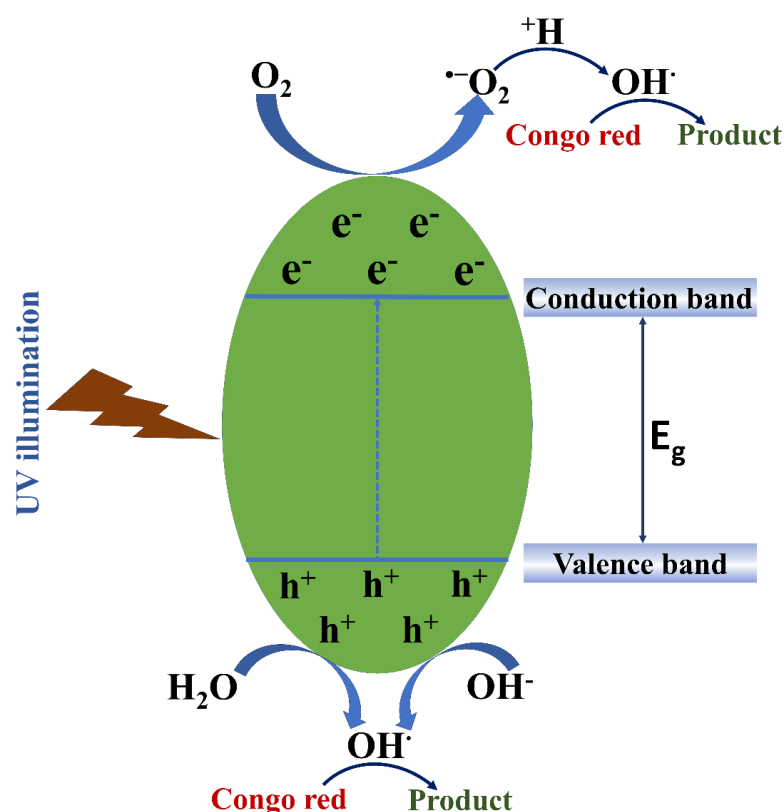


**Figure 4.7:** (a) Absorbance spectra of congo red (CR) dye added with exfoliated WS<sub>2</sub> until ~75 min exposure. Comparative plot for exfoliated, UV- and  $\beta$ -irradiated WS<sub>2</sub> as a photocatalyst added to CR dye for (b) 15 min, and (c) 45 min, respectively, under UV light exposure.

**Table 4.2.** Estimation of degradation percentage (%) of congo red dye calculated from the absorbance peak maxima.

Sl. No.	Sample	Degradation (%)	
		15 min illumination	45 min illumination
1	Exfoliated WS <sub>2</sub>	07.3	08.2
2	UV-irradiated WS <sub>2</sub>	07.3	10.2
3	$\beta$ -irradiated WS <sub>2</sub>	58.6	59.4





**Figure 4.8:** Schematic representation of a proposed mechanism of photocatalytic degradation of CR dye with WS<sub>2</sub> nanosystem as a photocatalyst under UV light illumination.

#### 4.8 Concluding remarks

In conclusion, the influence of UV- and  $\beta$ -irradiations on the morphological, vibrational, porosity, and dye degradation activity of the WS<sub>2</sub> system is studied in this chapter. The layered morphology of the WS<sub>2</sub> sheets revealed from FE-SEM imaging envisages the separation of layers from stacks of bulk WS<sub>2</sub> structure. The elemental analysis further confirmed the introduction of sulfur vacancies substantially. Raman spectroscopy was employed to investigate exfoliation in WS<sub>2</sub>, with the noticeable  $E'_{2g}$  and  $A_{1g}$  vibrational modes, located at  $\sim 352\text{ cm}^{-1}$  and  $\sim 421\text{ cm}^{-1}$ , respectively, exhibiting redshifts and broadening due to local heating effects with increasing laser power ( $\sim 2.5\text{ mW}$ ). Additionally, an observable increase in the  $E'_{2g}$ -to- $A_{1g}$  intensity ratio in UV-illuminated samples indicates improved exfoliation features under UV exposure. The enrichment of meso-macropores, along with structural defects in the WS<sub>2</sub> nanosheets, is attributed to the exfoliation process and the influence of radiation exposure on the system. Furthermore, the N<sub>2</sub> adsorption-desorption isotherm curve ( $\sim 77\text{ K}$ ) displayed a type-IV characteristic,



confirming the presence of mesopores, while pore-size distribution analysis based on the BJH model revealed the existence of various types of pores in the exfoliated WS<sub>2</sub> as well as after irradiation. A noticeable modification in the distribution of meso-macro pores was observed with radiation. Finally, the photocatalytic degradation efficiency was evaluated using congo red dye, with  $\beta$ -irradiated WS<sub>2</sub> demonstrating an enhanced degradation rate of approx. 59.38% upon UV illumination, for a time duration of 45 min.

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